

# ブラックホール磁気圏での 磁気リコネクションの数値計算

熊本大学

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ターサービス(株)

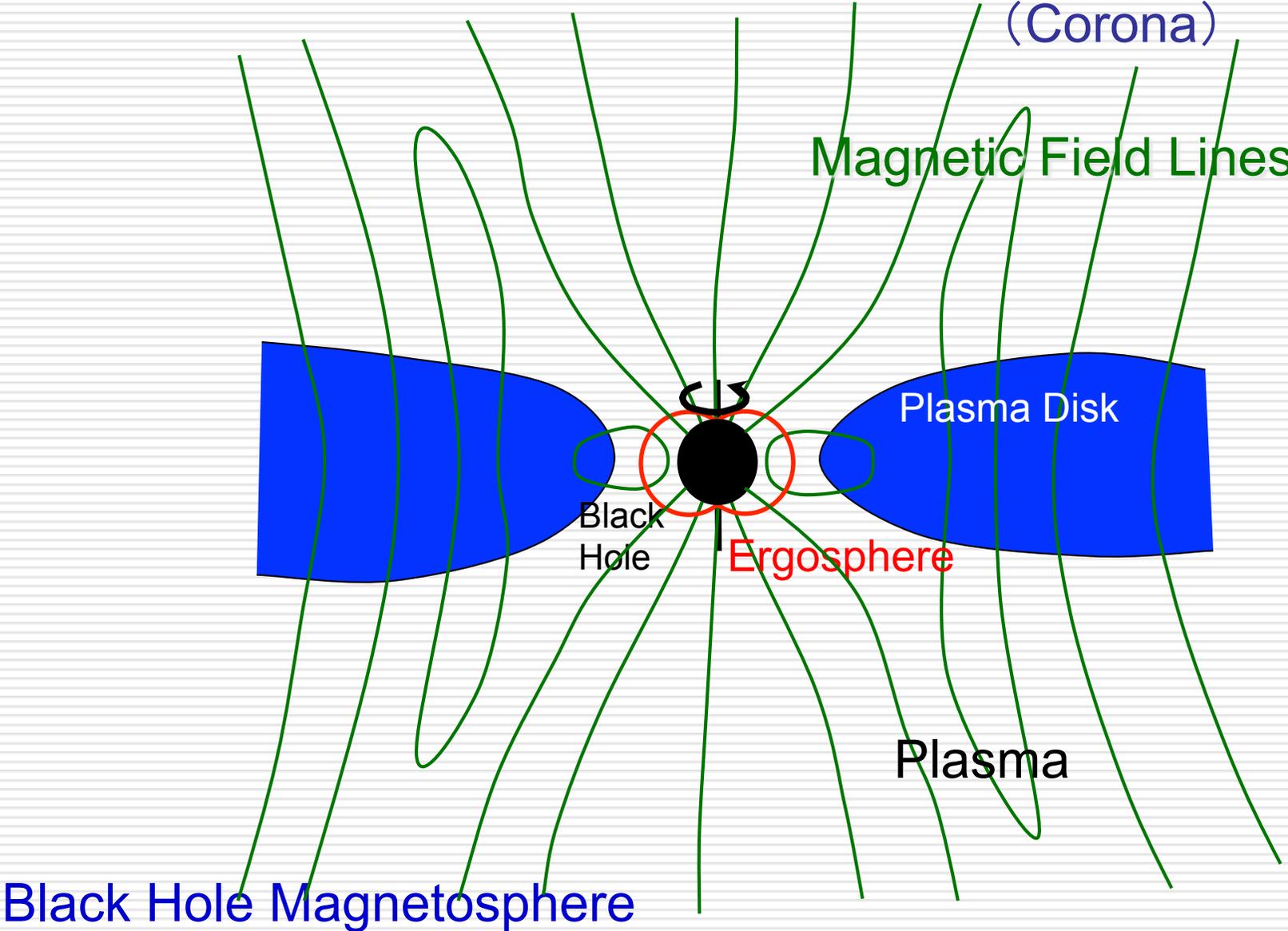
森野了悟

BHmag2012,名古屋大学, 2012.2.29

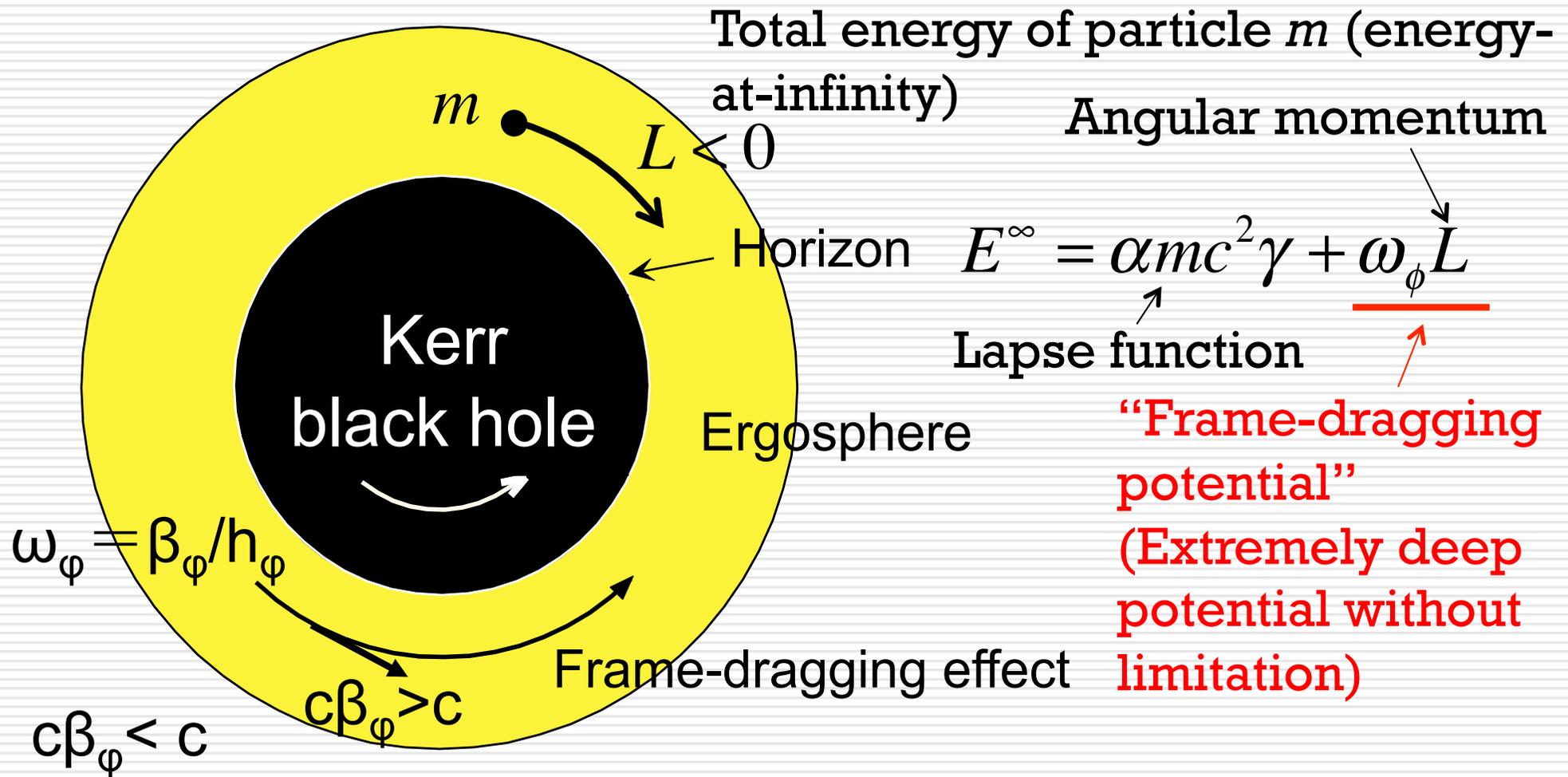
# Outline

- Motivation and basis: Magnetic reconnection around astrophysical black holes
- Standard equations of resistive GRMHD
- Test calculations of resistive GRMHD
- A simulation of magnetic reconnection with initially uniform magnetic field around rotating black hole
- Summary and Future plans

# Black Hole Magnetosphere

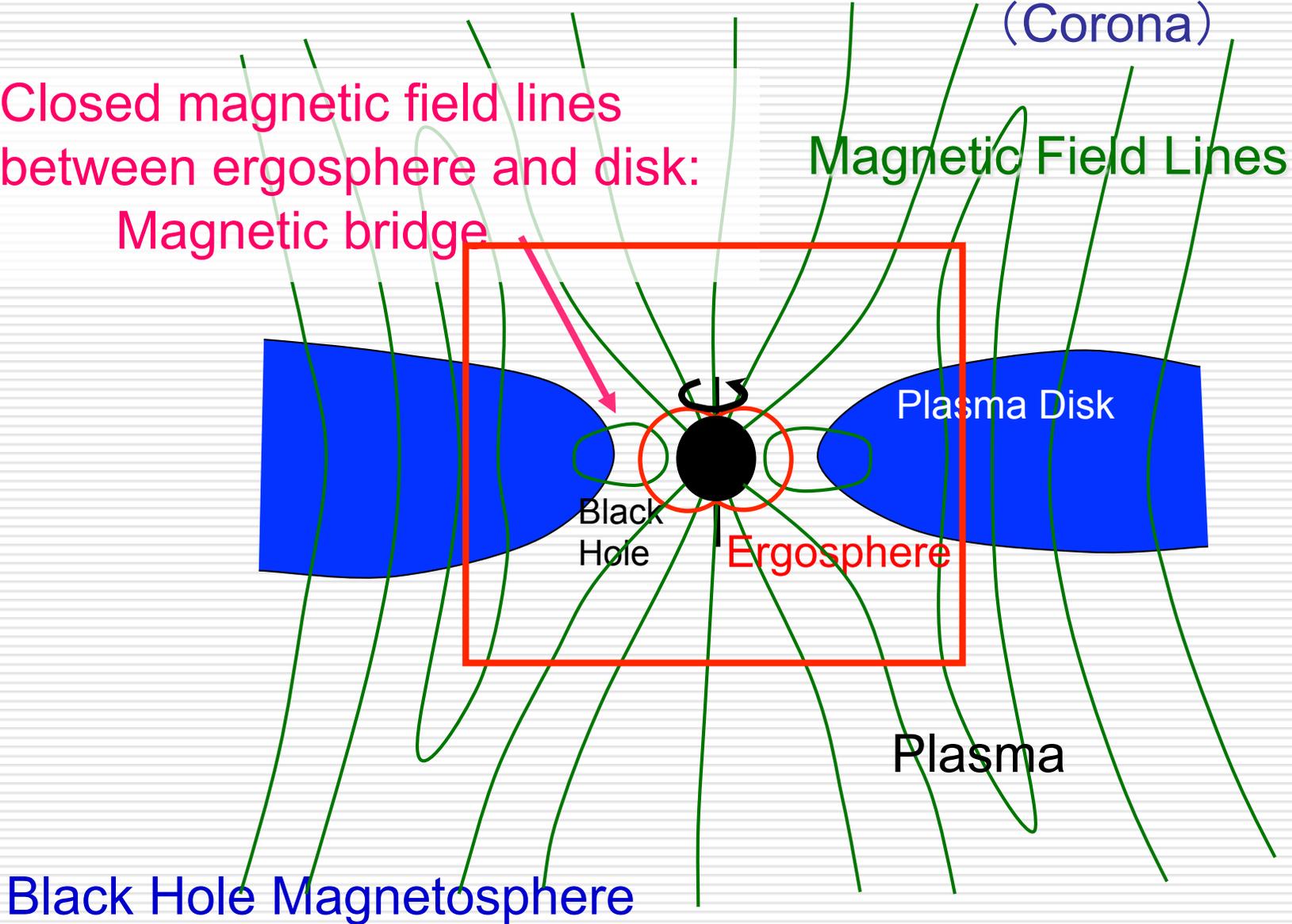


# Frame-dragging effect and ergosphere

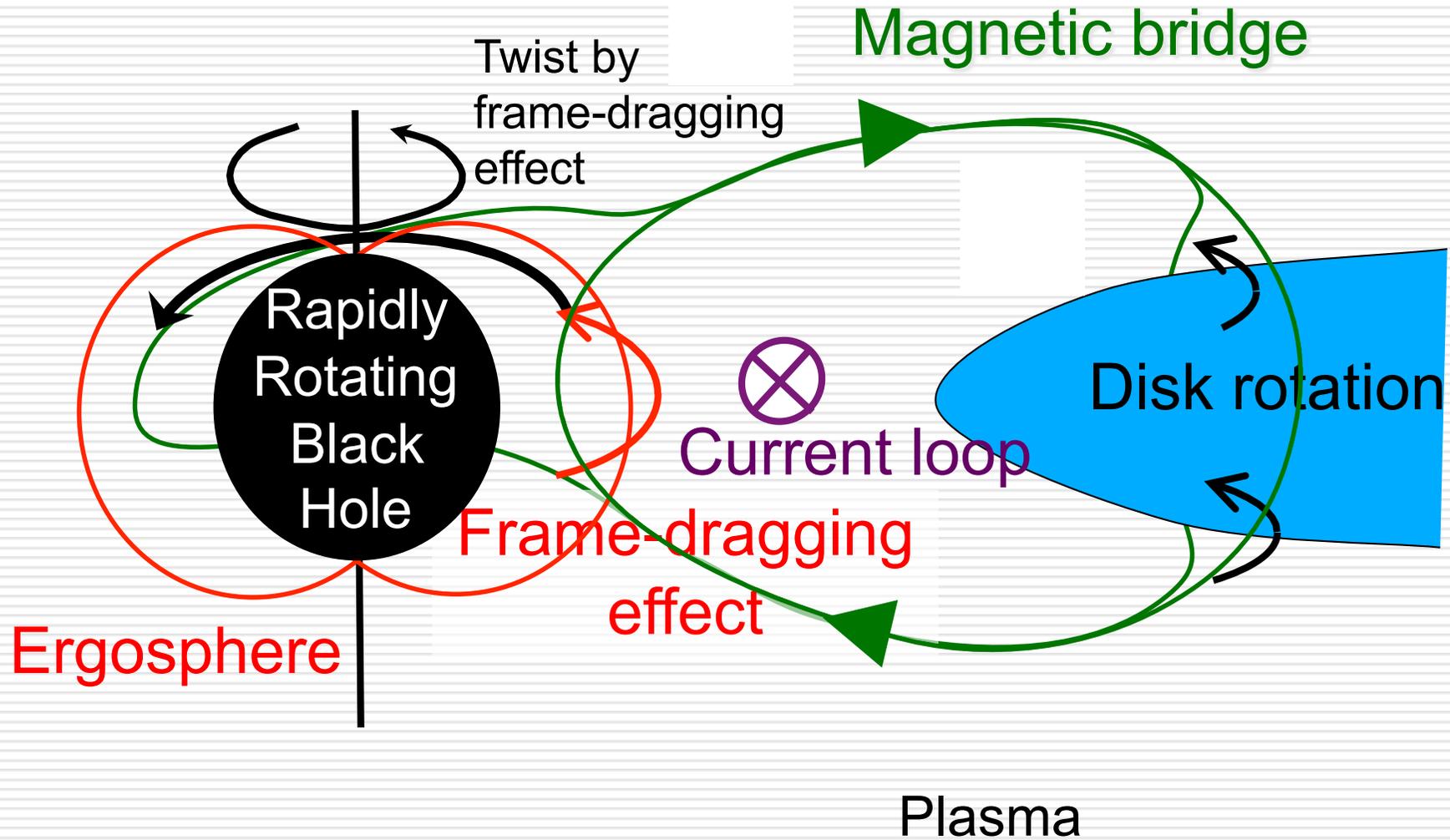


In the ergosphere, any matter, energy, and information can't move or propagate in the opposite direction of the black hole rotation. Plasma in the ergosphere behaves like very heavy plasma disk!

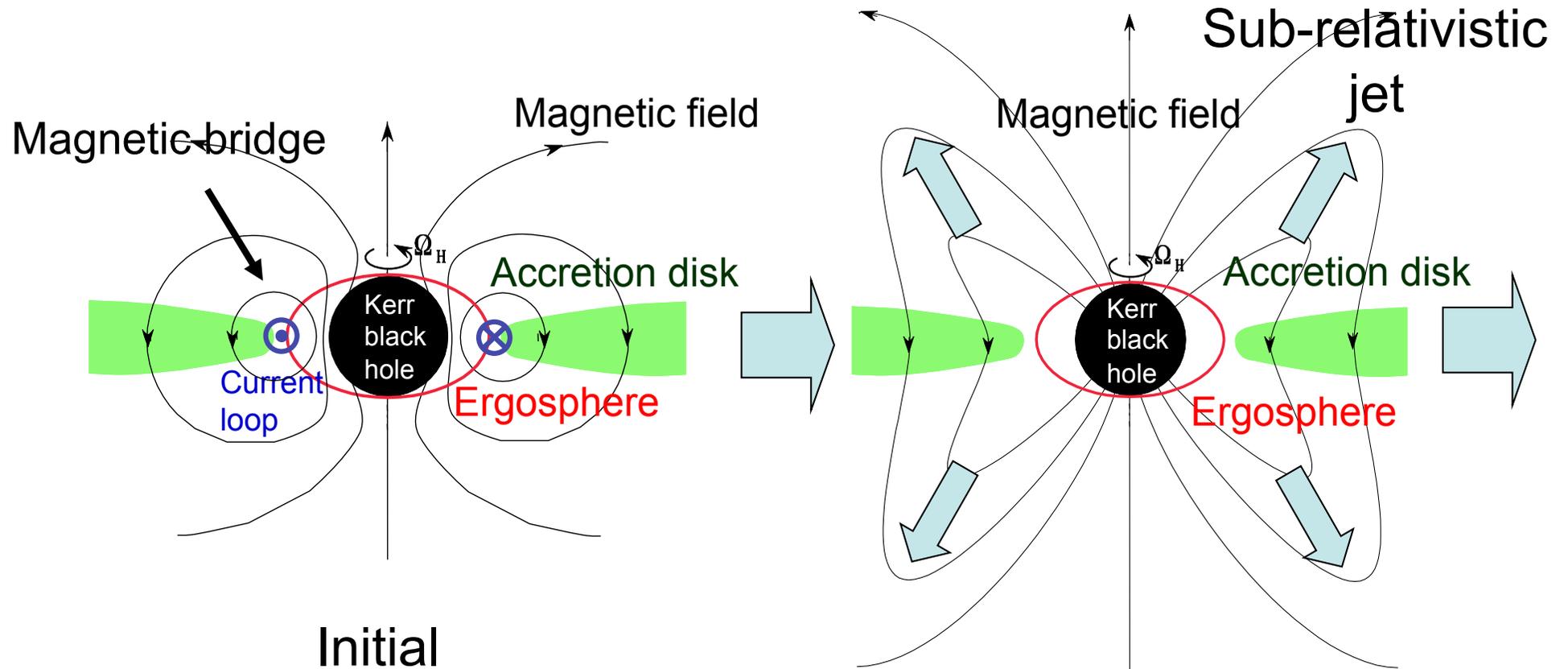
# Black Hole Magnetosphere



# Twist of magnetic bridge by ergosphere

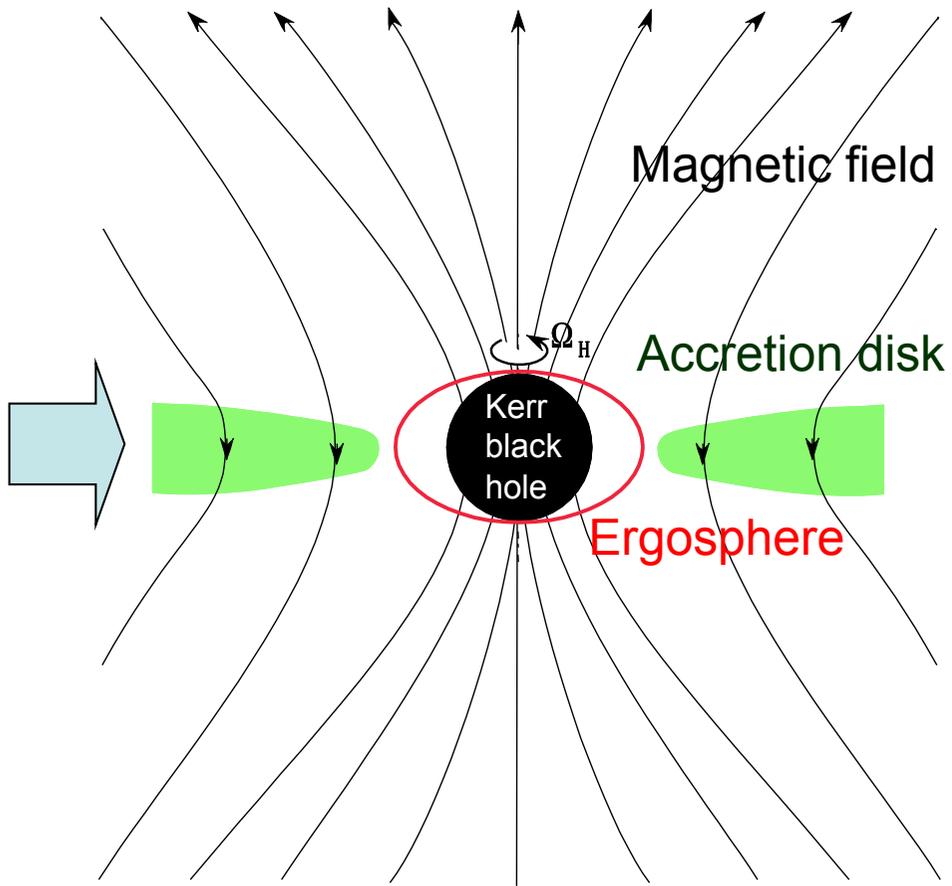


# Ideal GRMHD simulations: A phenomena caused by the magnetic bridge near the black hole

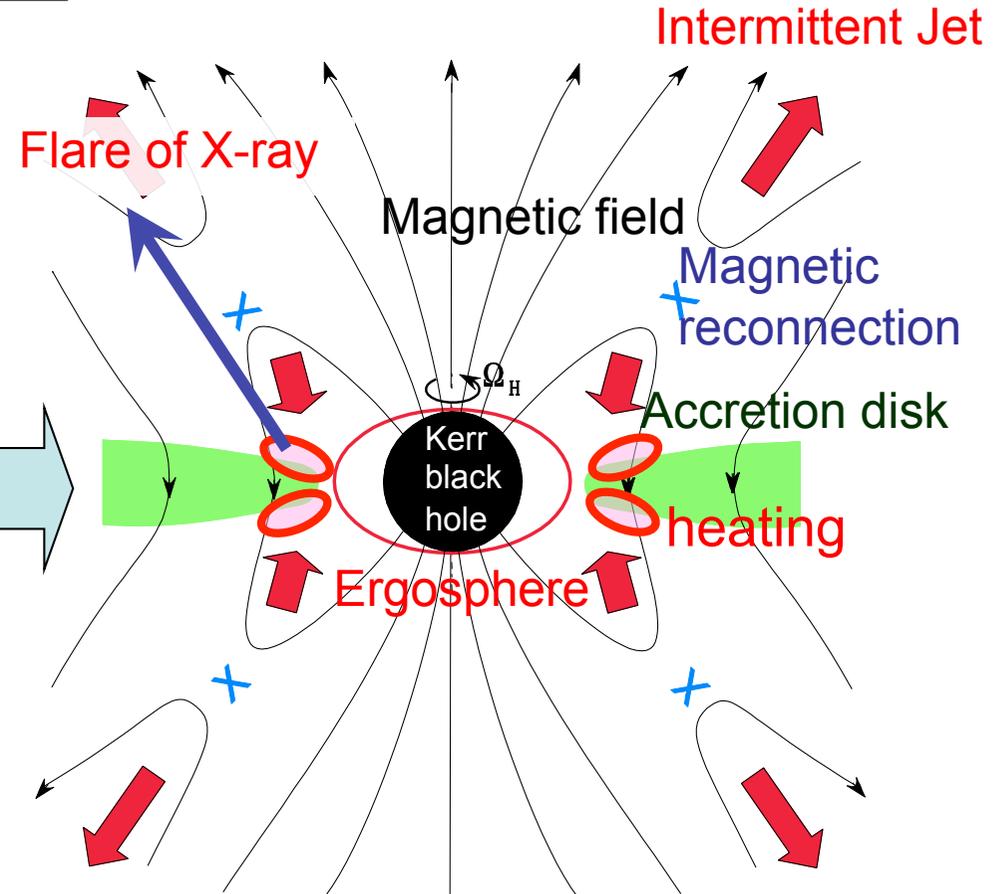


← Ideal GRMHD result

→ With finite resistivity



Anti-parallel magnetic field is formed



Magnetic reconnection must be important near black hole horizon!

GRMHD with finite resistivity

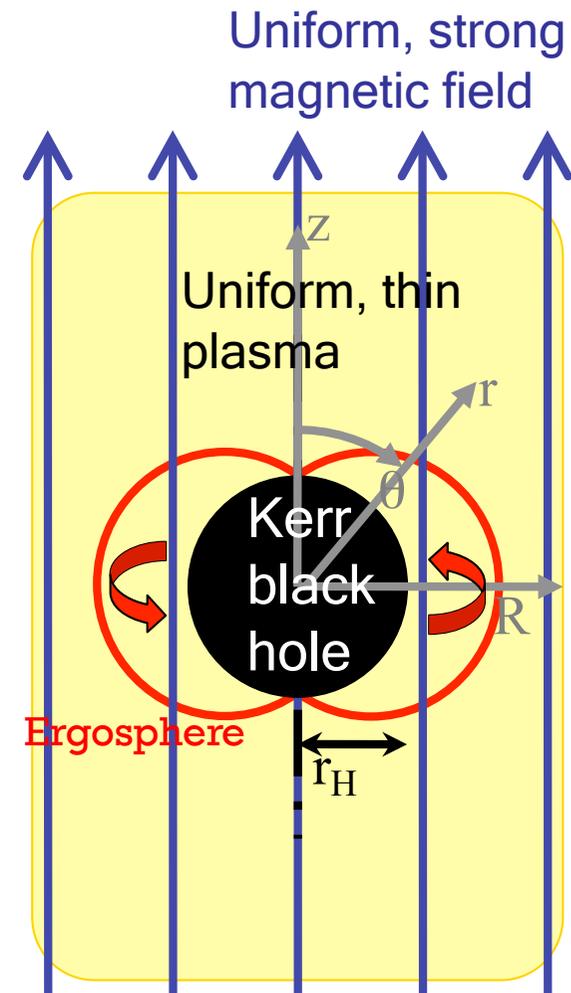
→ Near future important subject

# Even with uniform magnetic field as its initial condition . . .

Uniform magnetic field around rotating black hole with No Accretion disk :  
Initial condition

- (1) Kerr black hole: maximally rotating rotation parameter,  $a=J/J_{\max}=0.99995$
- (2) Magnetic field: Uniform around Kerr black hole (Wald solution)
- (3) Plasma: zero momentum, uniform, low density and pressure  
 $\rho_0=0.1B_0^2/c^2$ ,  $p_0=0.06\rho_0c^2$

(Koide, Shibata, Kudoh, Meier 2002)



A result of ideal GRMHD simulation

Power Radiation along Magnetic Field Lines cross Ergosphere,

Kerr black hole

Ergosphere

Magnetic field lines

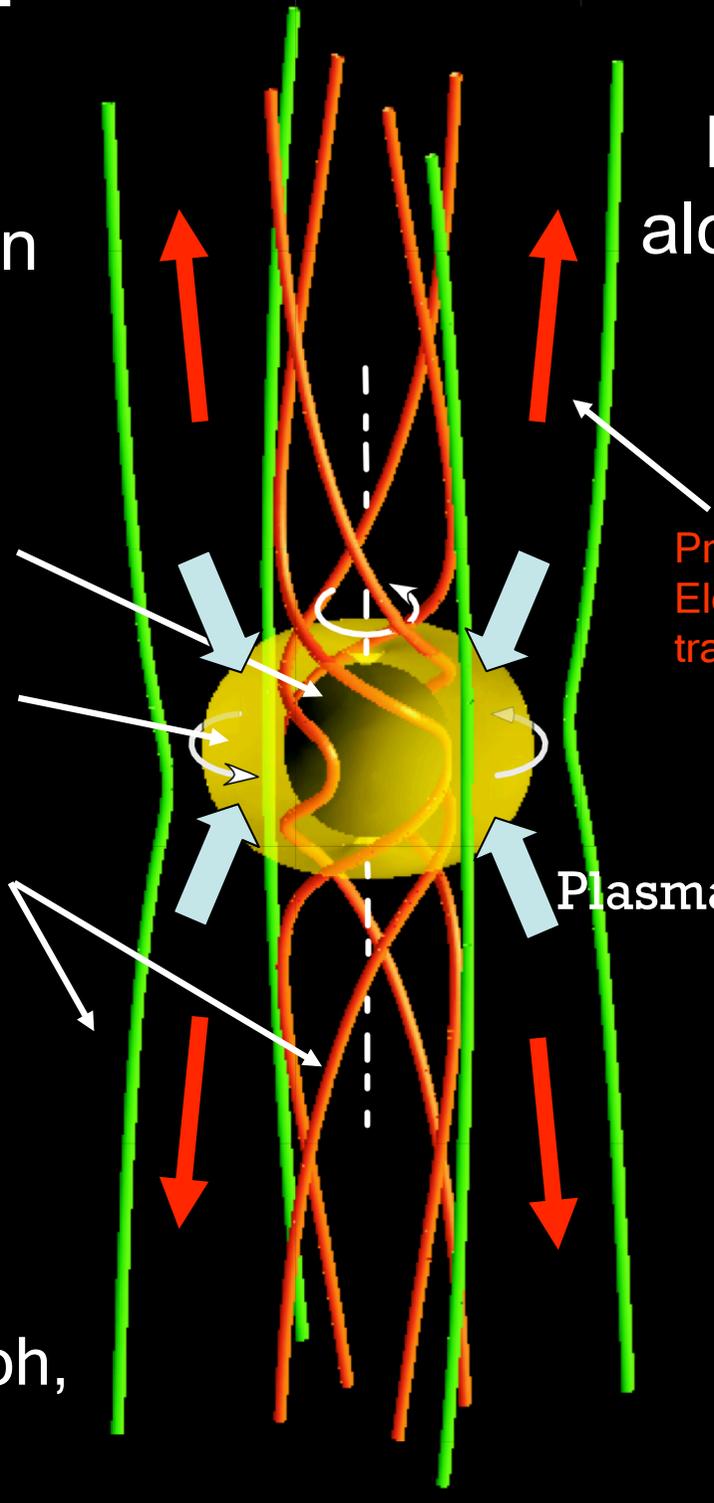
Plasma

Propagation of Alfvén wave:  
Electromagnetic energy transportation

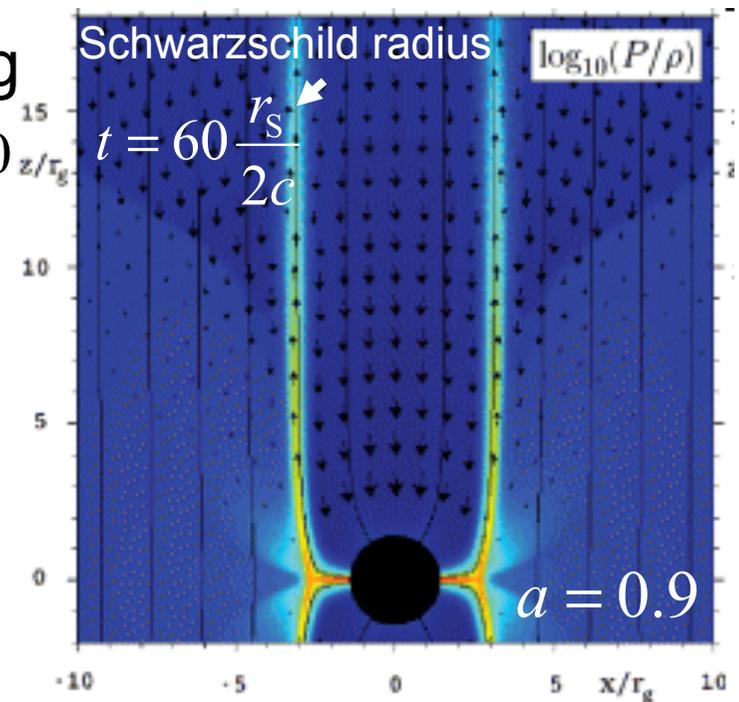
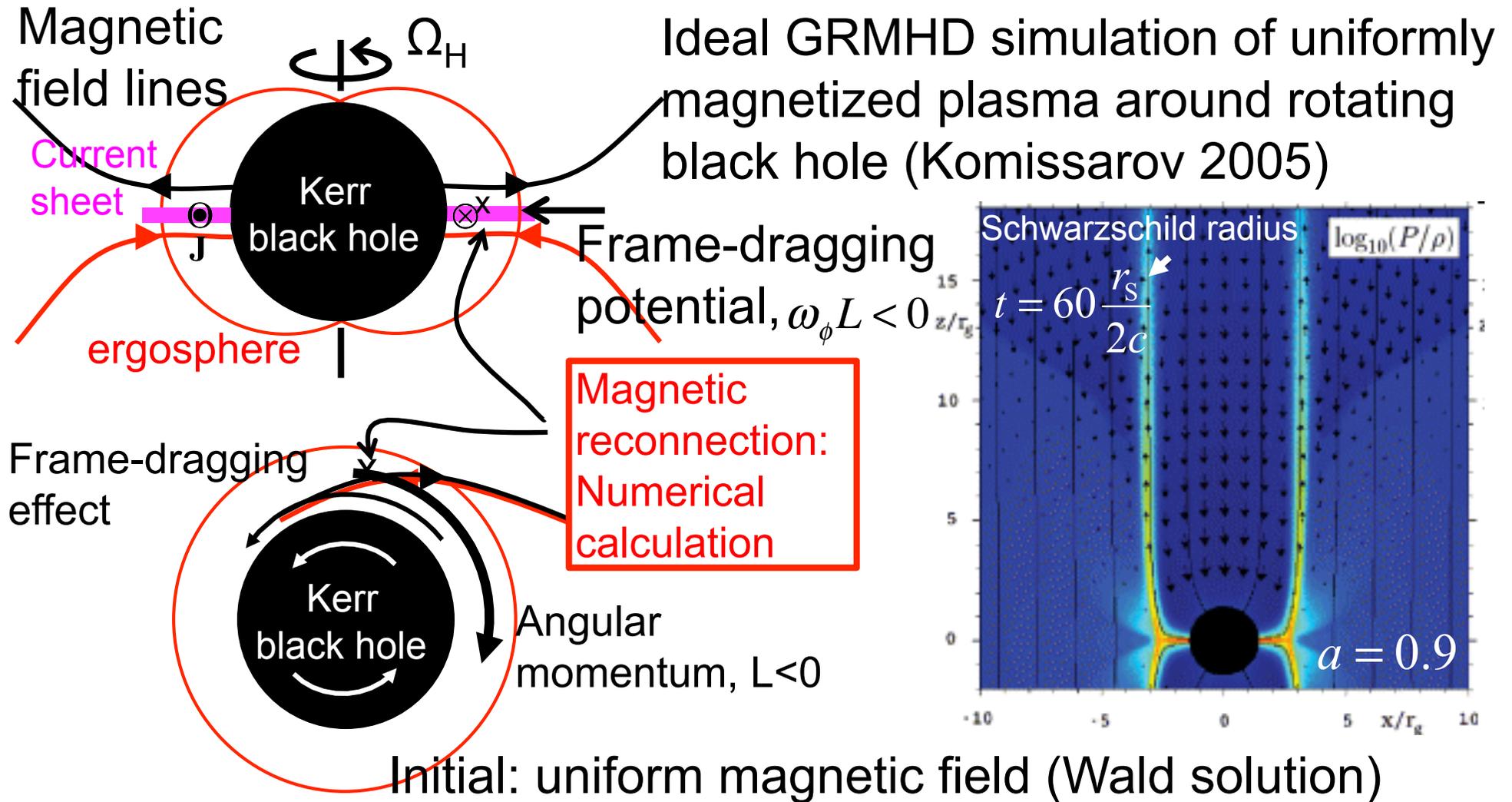
but No Outflow

$$t = 7\tau_S$$

Koide, Shibata, Kudoh,  
Meier 2002



# Longer term simulation of uniformly magnetized plasmas around Kerr BH



- ⇒ sprit monopole-like magnetic field in ergosphere
- ⇒ **magnetic reconnection** in ergosphere

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- **Standard equations of resistive GRMHD**
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# Covariant form of standard resistive GRMHD equations

Unit system  $\left[ \begin{array}{l} c = 1 \\ \mu_0 = 1 \end{array} \right.$

- General relativistic equations of conservation laws:

proper particle number density  
 $\nabla_\nu (n U^\nu) = 0$  (particle number)  
 4-velocity

Energy-momentum tensor  
 $\nabla_\nu T^{\mu\nu} = 0$  (energy and momentum)

Maxwell equations:

Field strength tensor  
 $\partial_\mu F_{\nu\lambda} + \partial_\nu F_{\lambda\mu} + \partial_\lambda F_{\mu\nu} = 0$   
 4-current density  
 $\nabla_\nu F^{\mu\nu} = -J^\mu$

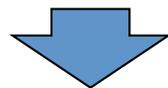
Ohm's law with resistivity:

resistivity  
 $F_{\mu\nu} U^\nu = \eta \left( J_\mu + (U_\nu J^\nu) U_\mu \right)$

Kerr Metric:

$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$   
 $g_{00} = -h_0^2; \quad g_{ii} = h_i^2; \quad g_{0i} = -h_i^2 \omega_i \quad (i = 1, 2, 3); \quad g_{ij} = 0 \quad (i \neq j)$

Lapse function:  $\alpha = \sqrt{h_0^2 + \sum_i (h_i \omega_i)^2}$



(gravitational time delay)

Shift vector:  $\beta_i = h_i \omega_i / \alpha$



$\beta = (\beta_1, \beta_2, \beta_3)$

(velocity of dragged frame)

# 3+1 form of resistive GRMHD equations for numerical calculations

$$\frac{\partial D}{\partial t} = -\nabla \cdot [\alpha D(\mathbf{v} + c\vec{\beta})] \quad (\text{conservation of particle number})$$

general relativistic effect

$$\frac{\partial \mathbf{P}}{\partial t} = -\nabla \cdot [\alpha(\mathbf{T} + c\vec{\beta}\mathbf{P})] - \left( D + \frac{\varepsilon}{c^2} \right) \nabla(c^2\alpha) + \alpha \mathbf{f}_{\text{curv}} - \mathbf{P} : \boldsymbol{\sigma}$$

(equation of motion)

special relativistic effect      Shear of frame-dragging

$$\frac{\partial \varepsilon}{\partial t} = -\nabla \cdot [\alpha(c^2\mathbf{P} - Dc^2\mathbf{v} + ec\vec{\beta})] - (\nabla\alpha) \cdot c^2\mathbf{P} - \mathbf{T} : \boldsymbol{\sigma} \quad (\text{equation of energy})$$

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times [\alpha(\mathbf{E} - c\vec{\beta} \times \mathbf{B})] \quad \alpha(\mathbf{J} + \rho_e c\vec{\beta}) + \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t} = \nabla \times \left[ \alpha \left( \mathbf{B} + \frac{\vec{\beta}}{c} \times \mathbf{E} \right) \right]$$

$$\nabla \cdot \mathbf{B} = 0 \quad \rho_e = \frac{\alpha}{c^2} \nabla \cdot \mathbf{E} \quad (\text{Maxwell equations})$$

$$\mathbf{E} + \mathbf{v} \times \mathbf{B} = \frac{\eta}{\gamma} \left[ \mathbf{J} - \gamma^2 \left( \rho_e - \frac{1}{c^2} (\mathbf{v} \cdot \mathbf{J}) \right) \mathbf{v} \right] \quad (\text{Ohm's law with finite resistivity})$$

All of these equations are coupled! We have to consider Ampere's law and Gauss's law of electric charge. We use the numerical method of Watanabe & Yokoyama (2006).

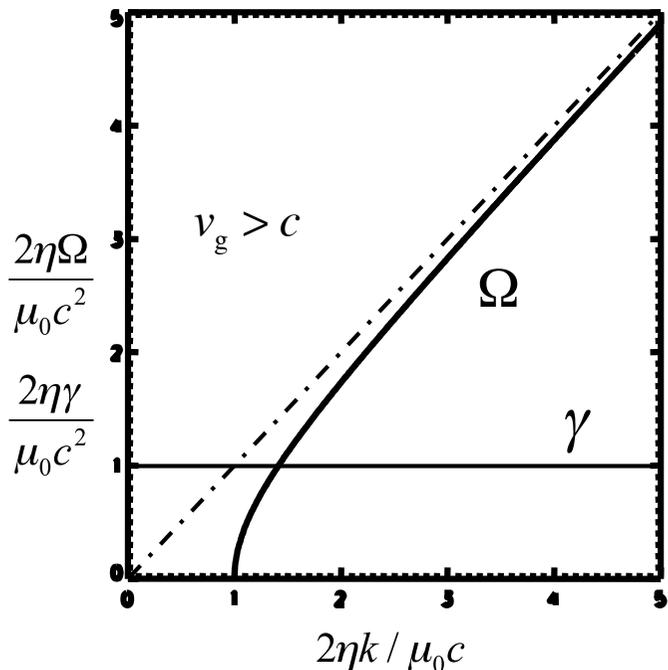
# “Standard” resistive GRMHD equations※

=Relativistic one-fluid model with resistivity

- Conservation of particle number, momentum, and energy
- Maxwell equations
- Standard relativistic Ohm’s law :

$$U^\nu F^\mu{}_\nu = \eta (J^\nu - \rho_e U^\nu)$$

Resistivity



Dispersion relation of electromagnetic wave in resistive plasma

Group velocity of electromagnetic wave is greater than light speed

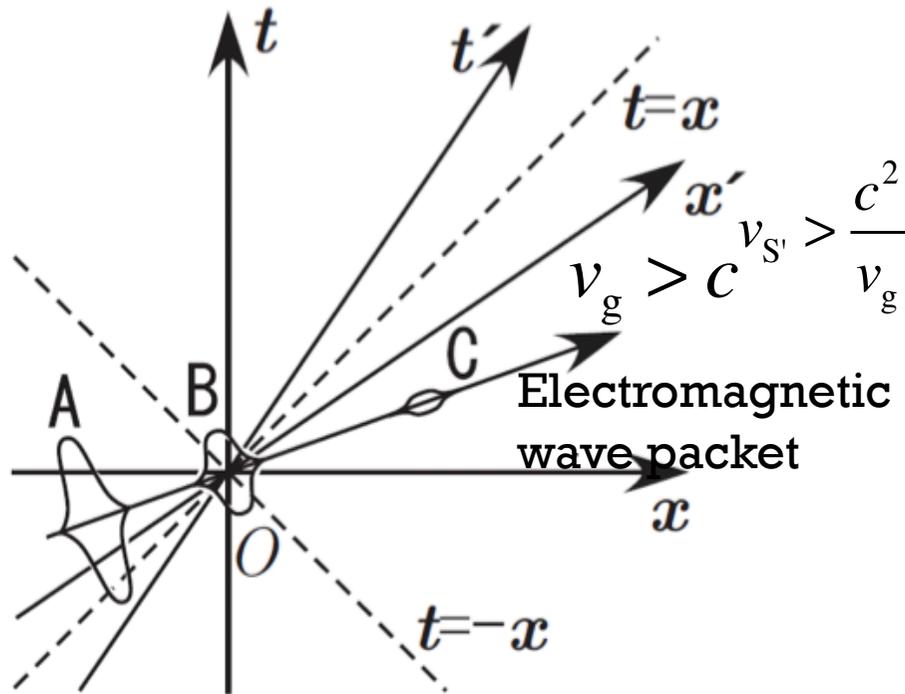


**Causality broken?**

※ Now, numerical calculations are limited in special relativistic version.

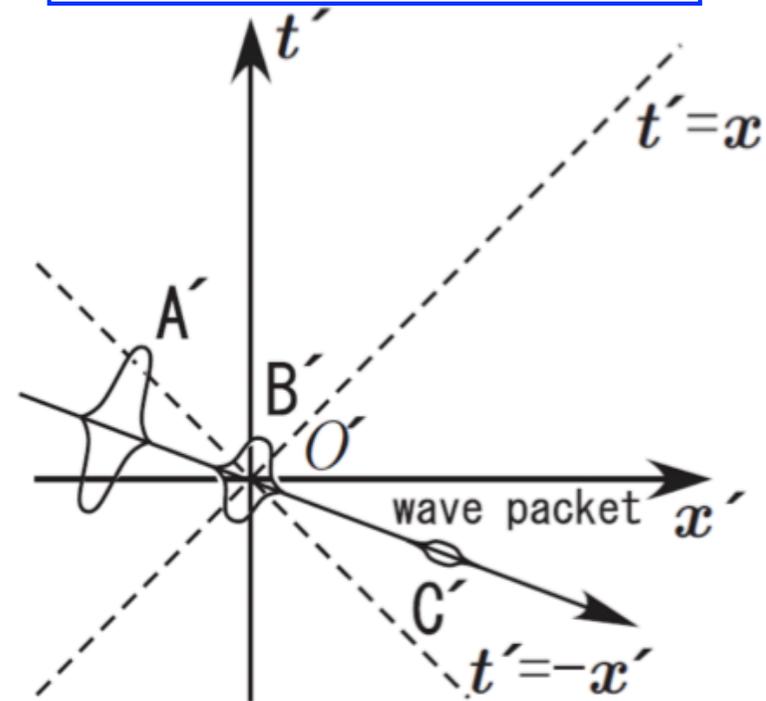
# Superluminal information transportation: Break down of causality?: Spontaneous decrease in entropy

Plasma rest frame



increase in entropy  
 electromagnetic energy  
 $\Rightarrow$  thermal energy

New coordinates frame



decrease in entropy  
 thermal energy  $\Rightarrow$   
 electromagnetic energy

# Generalized GRMHD equations

$$\nabla_\nu (\rho U^\nu) = 0$$

$$\nabla_\nu \sqrt{h} U^\nu = 0$$

Equation of 3+1 formalism required for numerical calculation of the generalized GRMHD are much complicated so that no one wants to perform calculation with them.



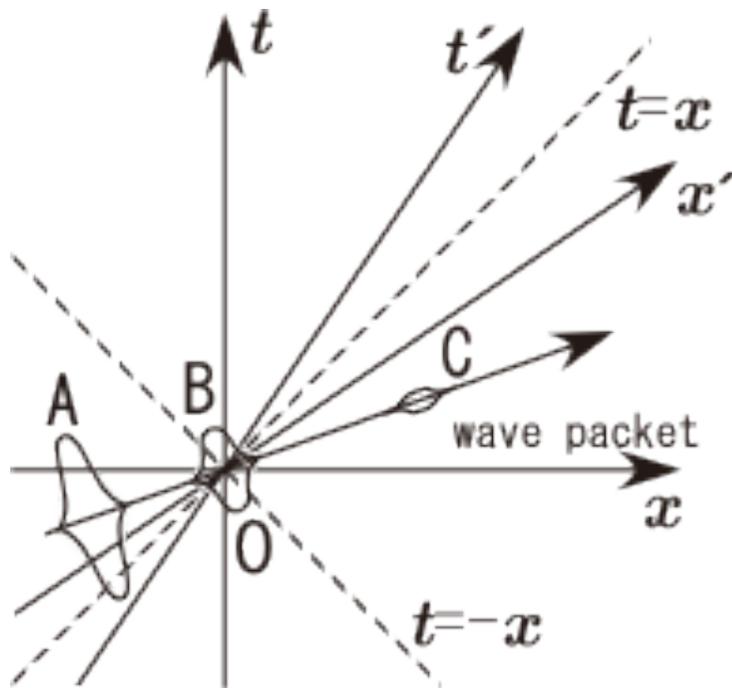
Therm To clarify the limitation of the standard resistive GRMHD equations, we tried to perform acausal phenomena expected (not welcome) in the standard equations.

$$\eta < \left( \sim 3 \right) \frac{m}{e\rho} \sqrt{\mu h^\dagger} \Rightarrow v_g \leq c$$

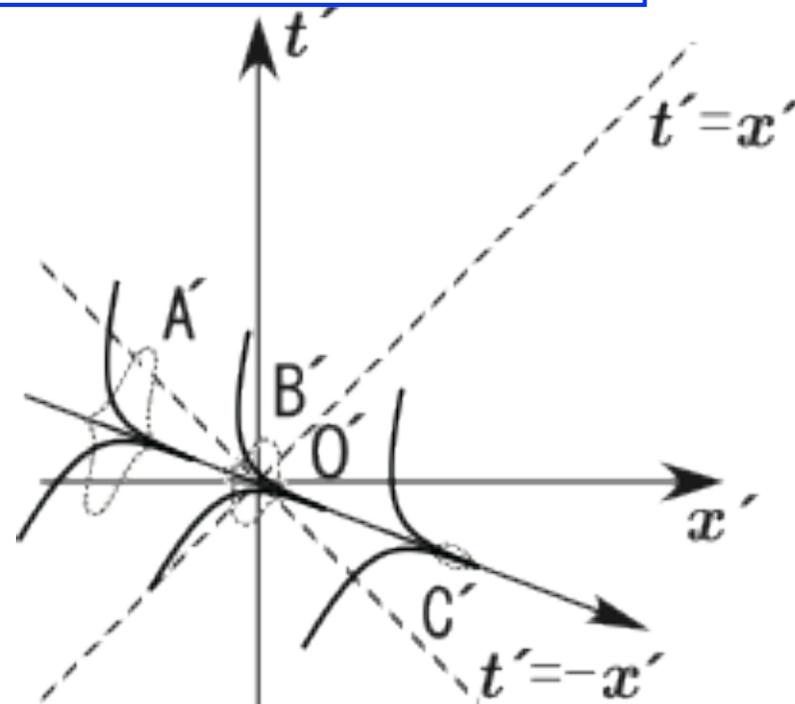
standard Ohm's law  
general relativistic Ohm's law

In the new coordinates with fast plasma flow, the wave packet has monotonic amplitude profile and then could not send a signal

Plasma rest frame



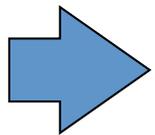
New coordinates frame



It seems that monotonic wave decreases. This shape of the wave is not “wave packet” anymore. The wave packet is not Lorentz invariant.

# Remarks on causality problem of RRMHD equations

- Wave-packet is not Lorentz invariant, and then cannot be accept as a relativistic object.
- The group velocity, which represents propagation velocity of wave packet, is not relativistic quantity. In fact, the definition of the group velocity,  $v_g = \partial\omega / \partial k$ , is not scalar nor covariant variable.
- We cannot discuss causality using the group velocity.
- Instead of group velocity, “head velocity”  $v_h = \lim_{k \rightarrow \infty} \omega / k$  presents the propagation velocity of information. In resistive RMHD, the head velocity is light speed, and Lorentz invariant.



**Concerning causality, there is no limitation of standard resistive GRMHD!**

- Koide & Morino, Phys. Rev. D, 84, 083009, Jackson, Classical Electrodynamics (1962).

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# Test calculations of resistive GRMHD

- Interaction between plasma and uniform magnetic field around black hole

# Plasma fall in uniform magnetic field into a Schwarzschild black hole

$$\eta = 10^{-2} r_s c$$

( $\sim$  ideal RMHD)

Color: pressure

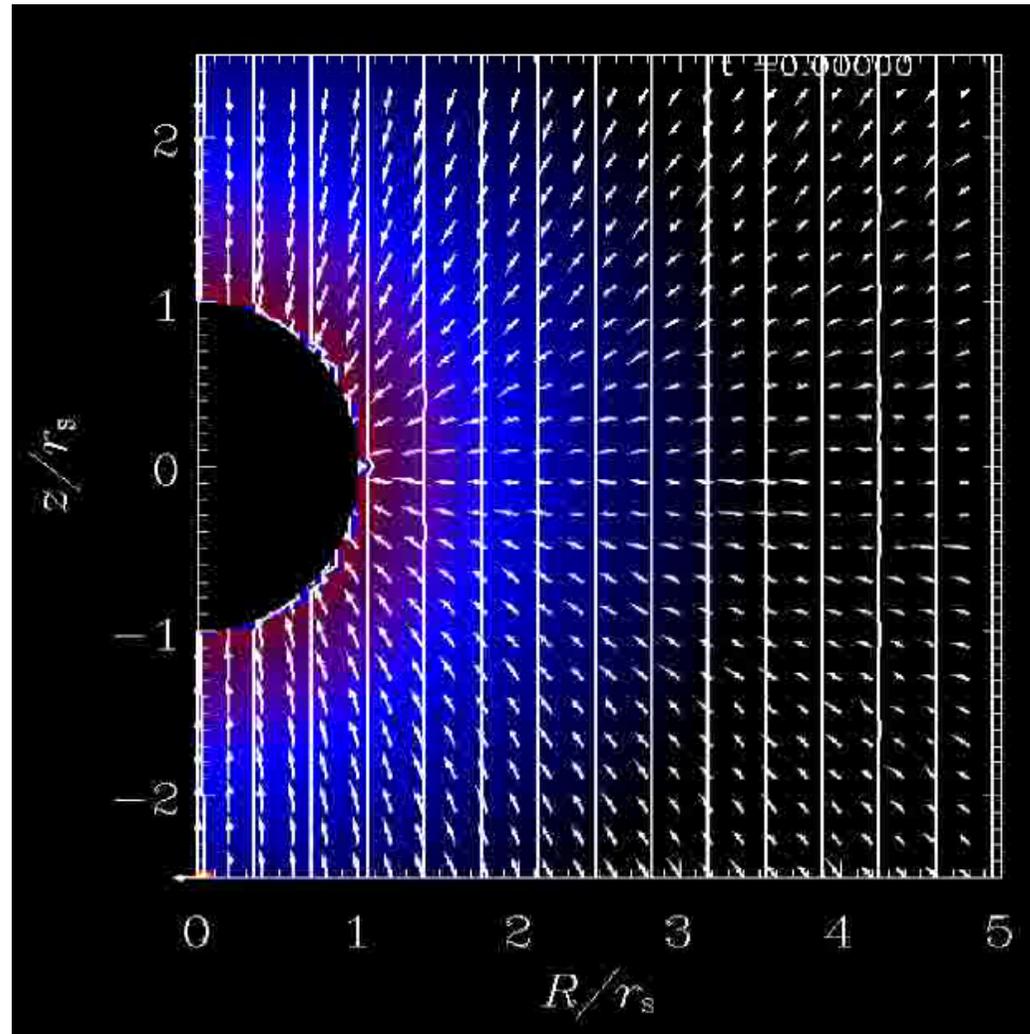
$$a = 0$$

Unit of length

$$r_s = \frac{2GM}{c^2}$$

Unit of time

$$\tau_s = \frac{r_s}{c}$$



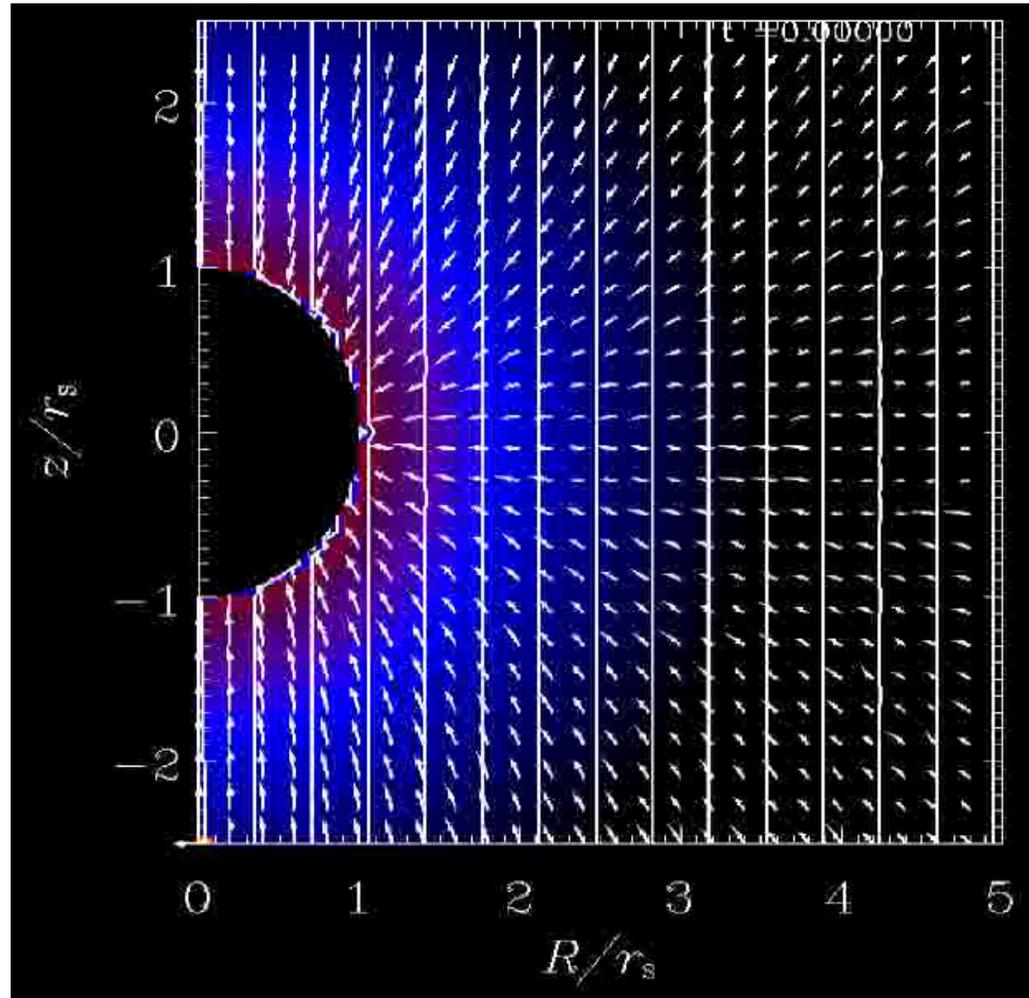
# Plasma fall in uniform magnetic field into a Schwarzschild black hole

$$\eta = 10^3 r_s c$$

( $\sim$  Vacuum)

Color: pressure

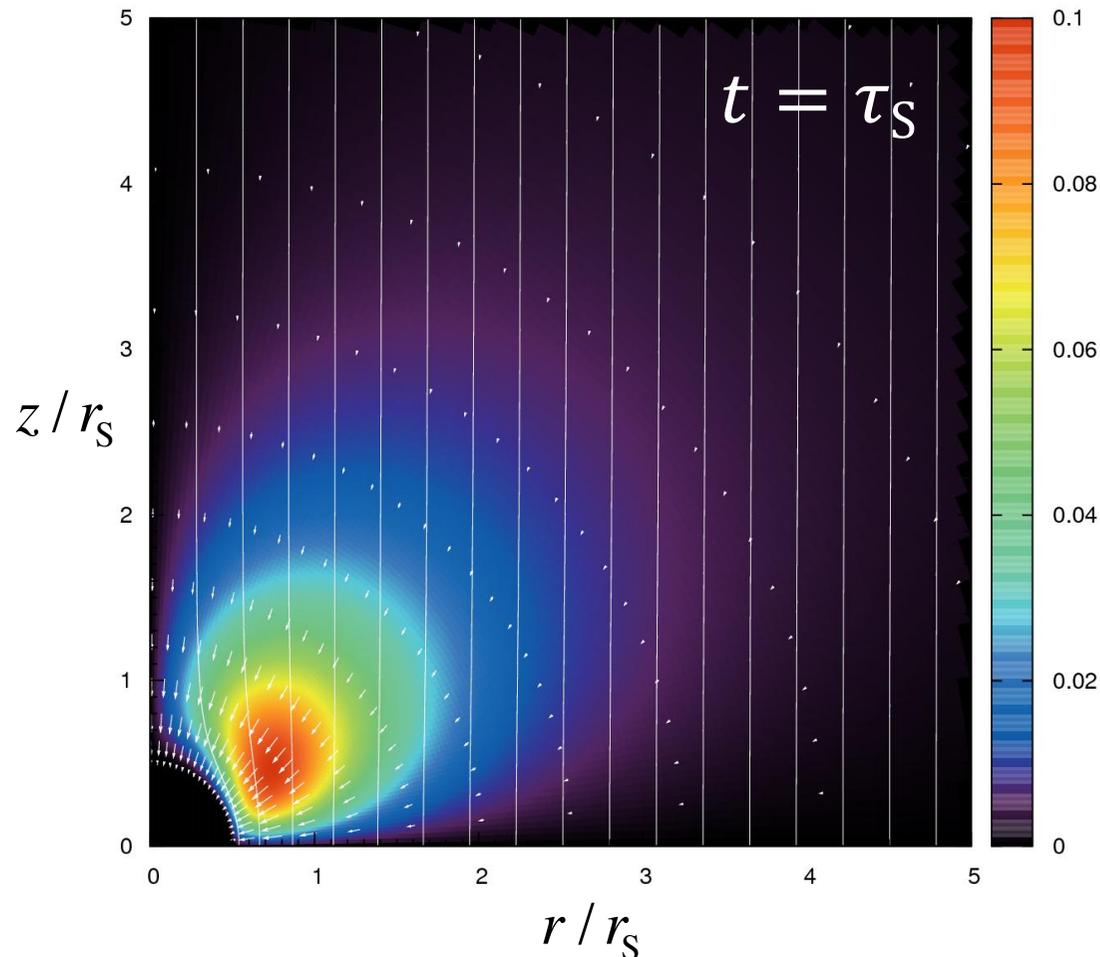
$$a = 0$$



# Test calculation: Frame-dragging dynamo effects

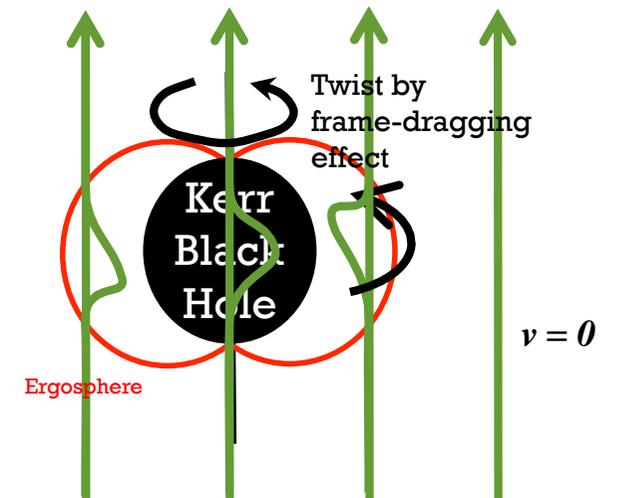
$$B_\phi / B_0$$

$$\eta = 10^{-3} r_S c$$



$$a \equiv \frac{J}{J_{\max}} = 1$$

$$J_{\max} = \frac{GM^2}{c}$$

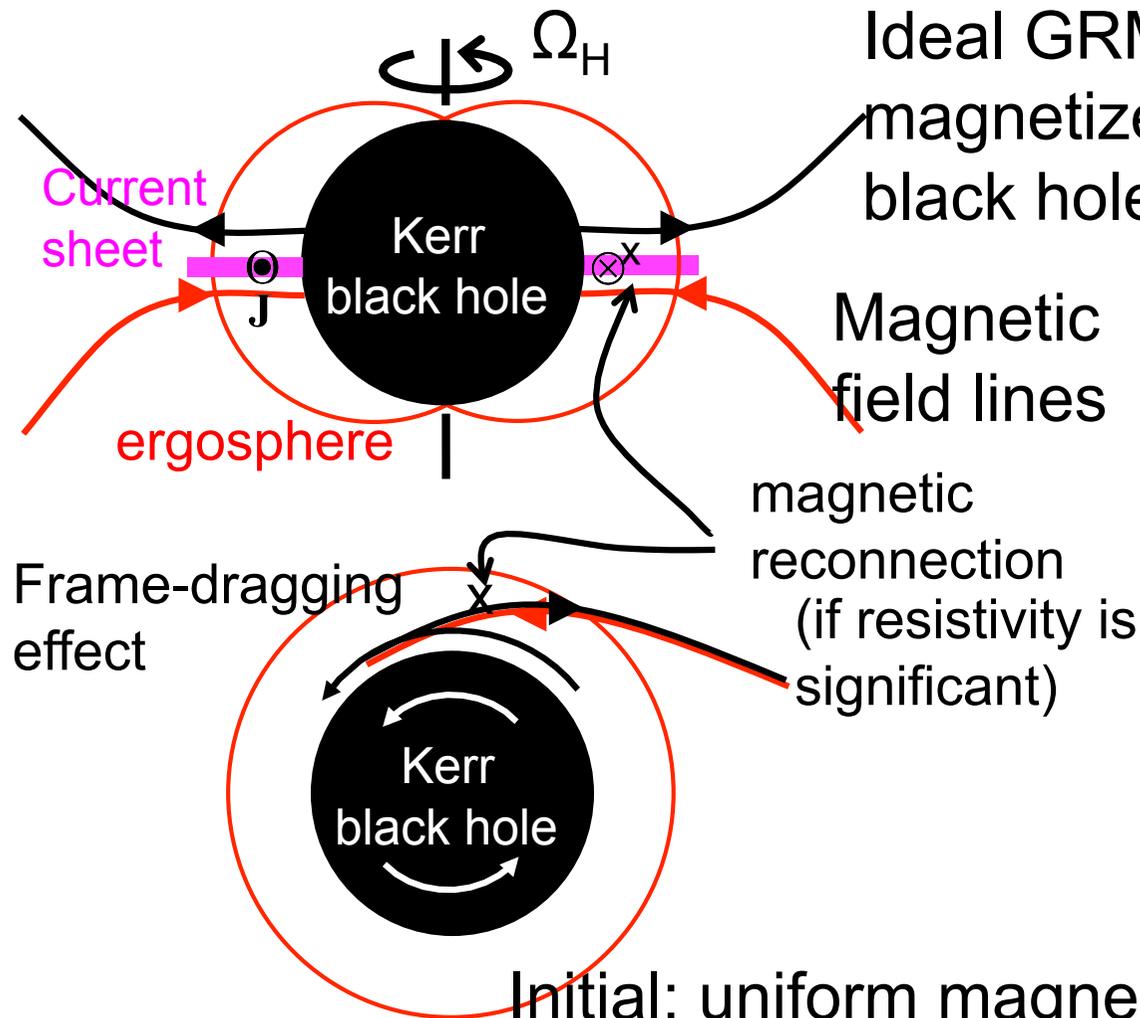


HLL+MUSCL scheme

# Outline

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- A simulation of magnetic reconnection with initially uniform magnetic field around rotating black hole
  - Resistive GRMHD simulation with initial uniform magnetic field around a rotating black hole (2 trial)
  - Resistive GRMHD simulation with initial split monopole magnetic field around a rotating black hole
- Summary and Future plans

# Longer term simulation of uniformly magnetized plasmas around Kerr BH

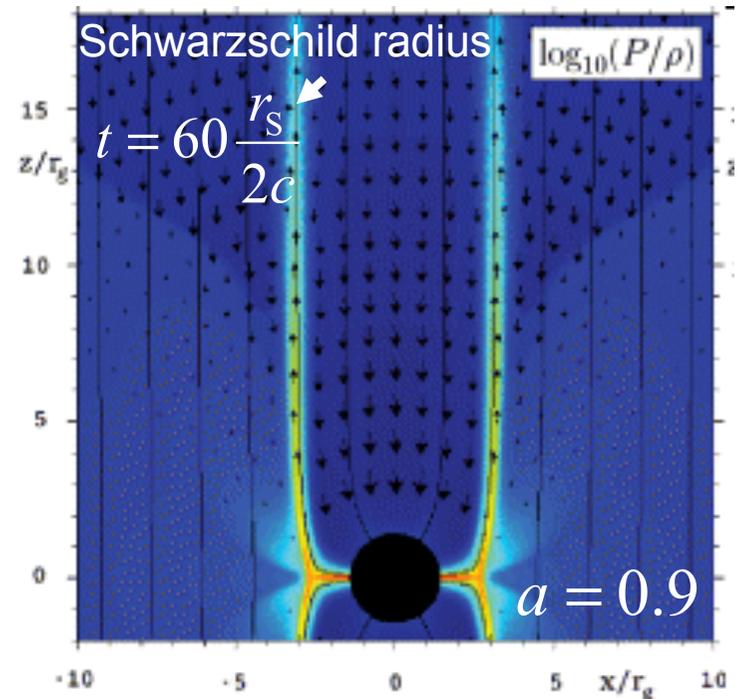


Ideal GRMHD simulation of uniformly magnetized plasma around rotating black hole (Komissarov 2005)

Magnetic field lines

magnetic reconnection (if resistivity is significant)

Frame-dragging effect



- Initial: uniform magnetic field (Wald solution)
- ⇒ split monopole-like magnetic field in ergosphere
- ⇒ **magnetic reconnection** in ergosphere

# Trial calculation with uniform magnetic field as initial field

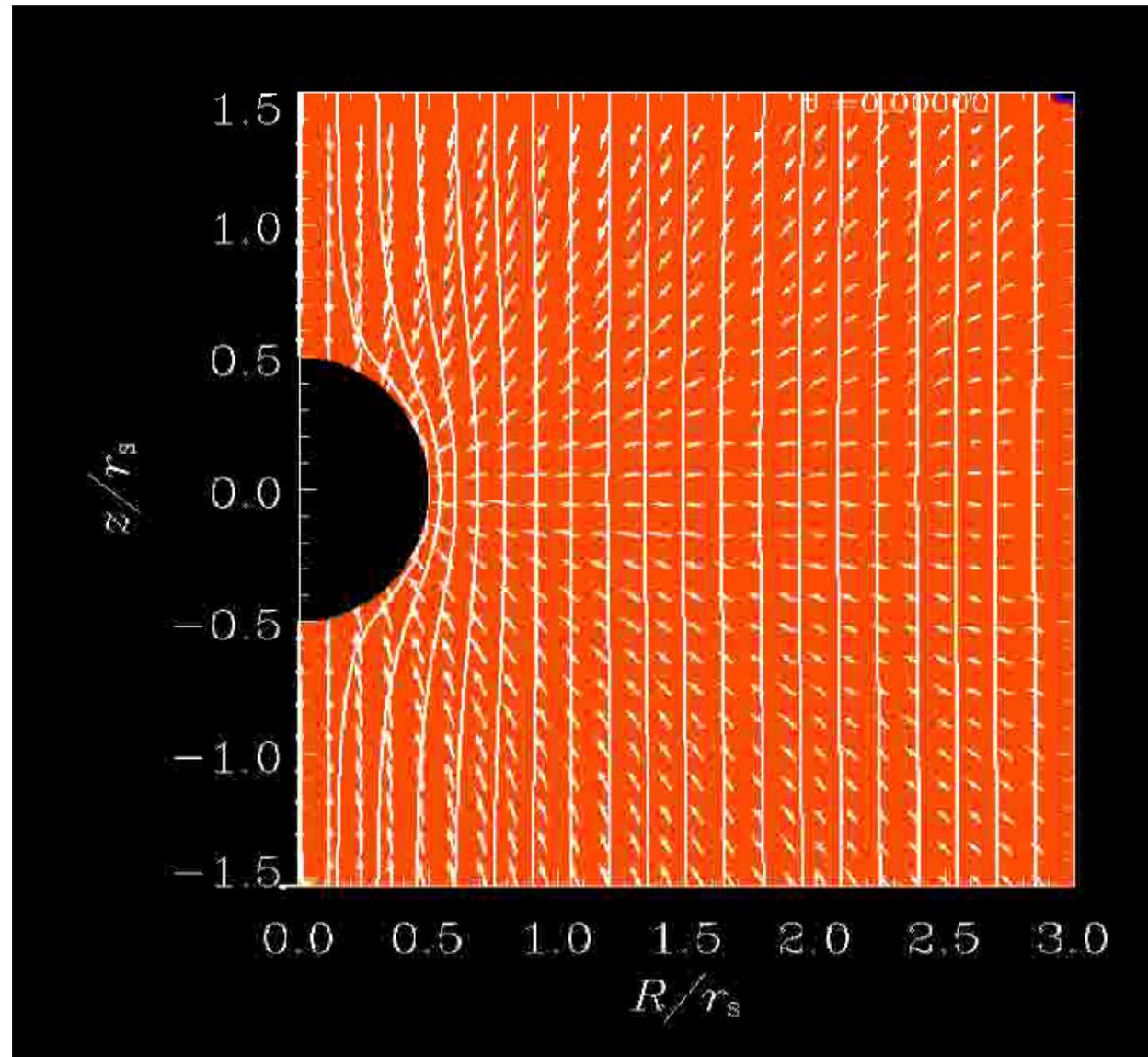
$$a = 0.999995$$

$$\beta = 0.025$$

Color:  $B_\phi / \rho^{1/2}$

Lines: magnetic field

Arrows: velocity



# Trial calculation with uniform magnetic field (try #2)

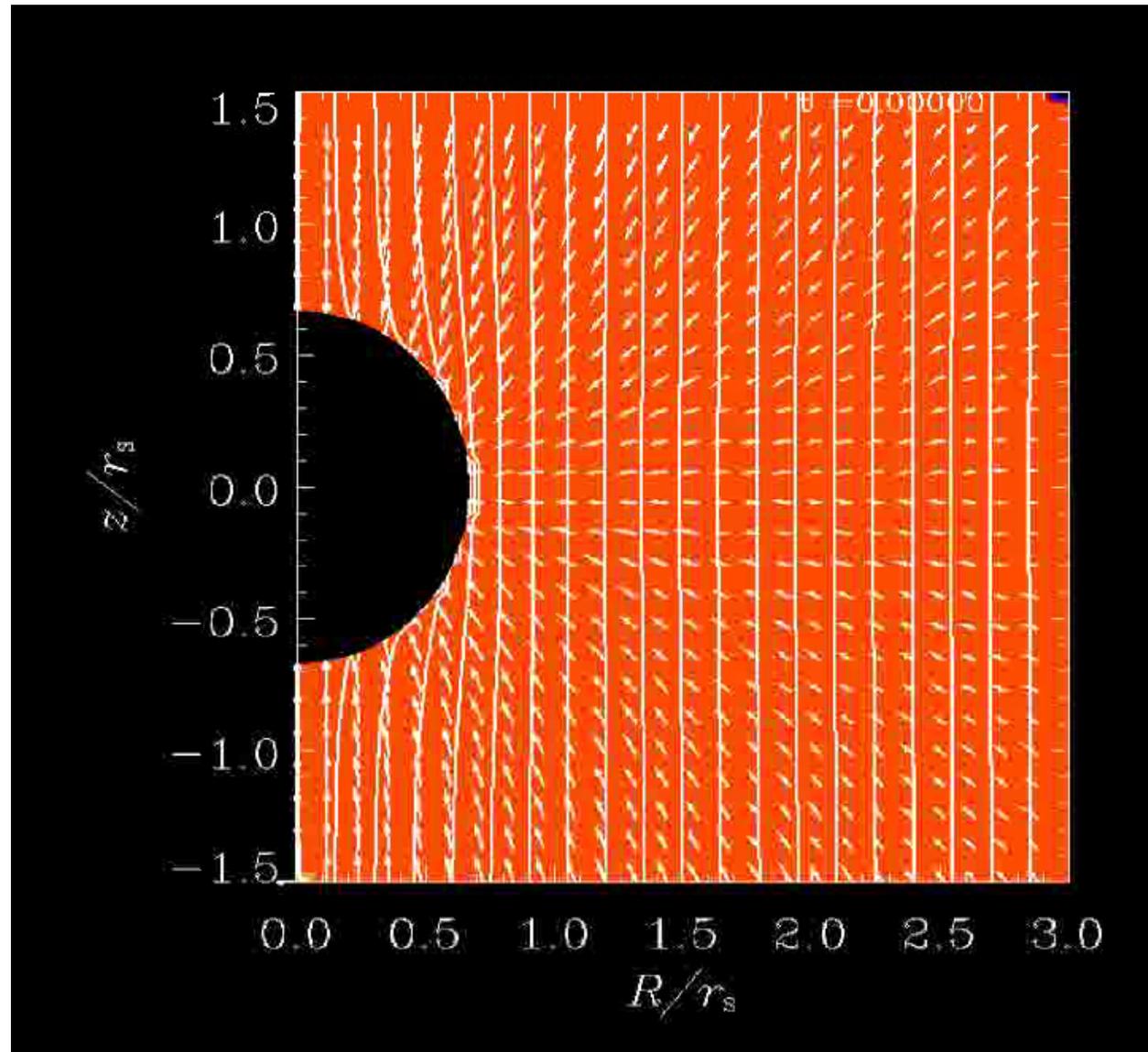
$$a = 0.94$$

$$\beta = 0.025$$

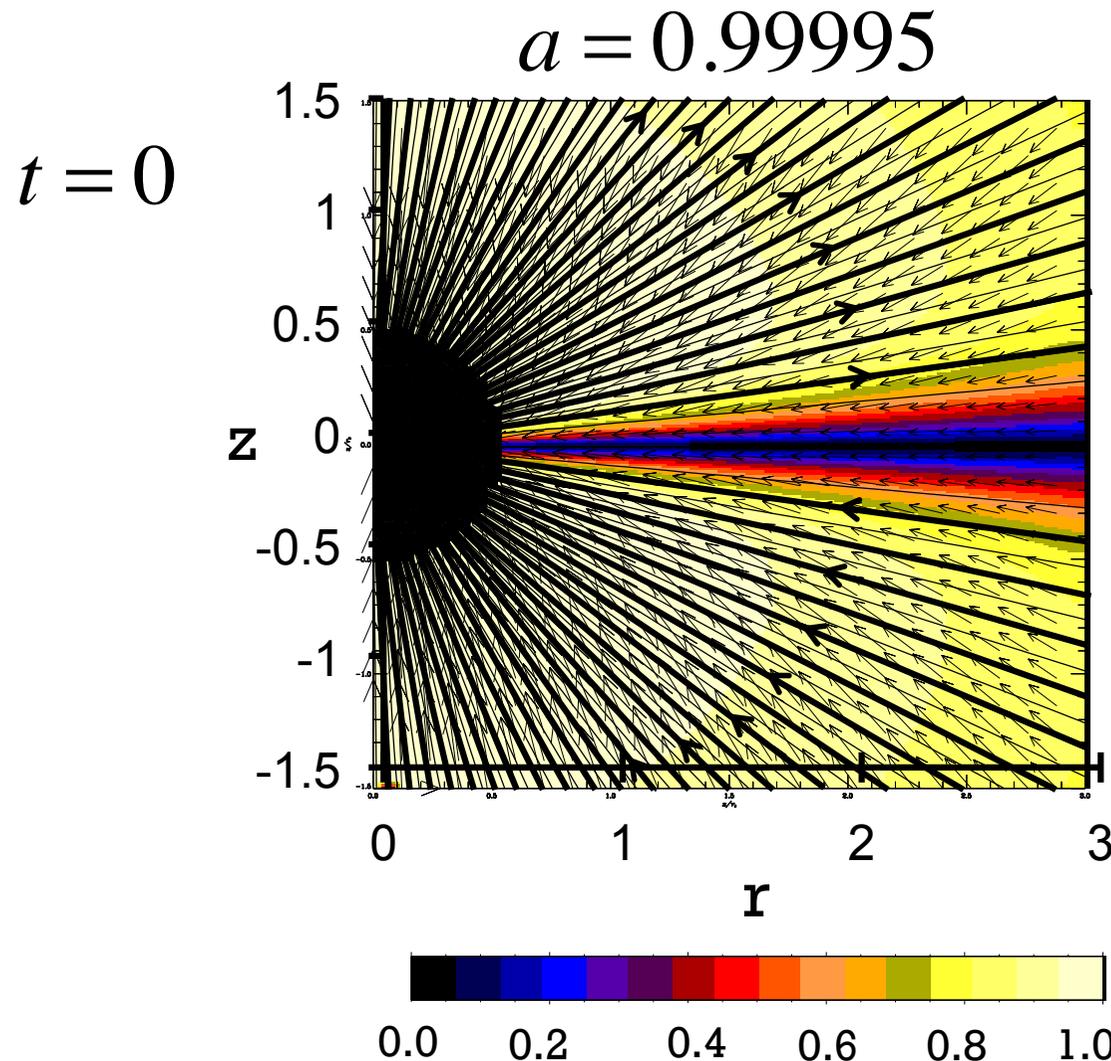
Color:  $B_\phi / \rho^{1/2}$

Lines: magnetic field

Arrows: velocity



# Initial condition of magnetic field of a simple case: Split monopole field



Resistivity

$$\eta = 0.01 r_s$$

Relativistic Alfvén velocity

$$v_A = \frac{B}{\sqrt{\rho c^2 + \frac{\Gamma}{\Gamma-1} p + B^2}} c$$

Split-monopole  
magnetic field

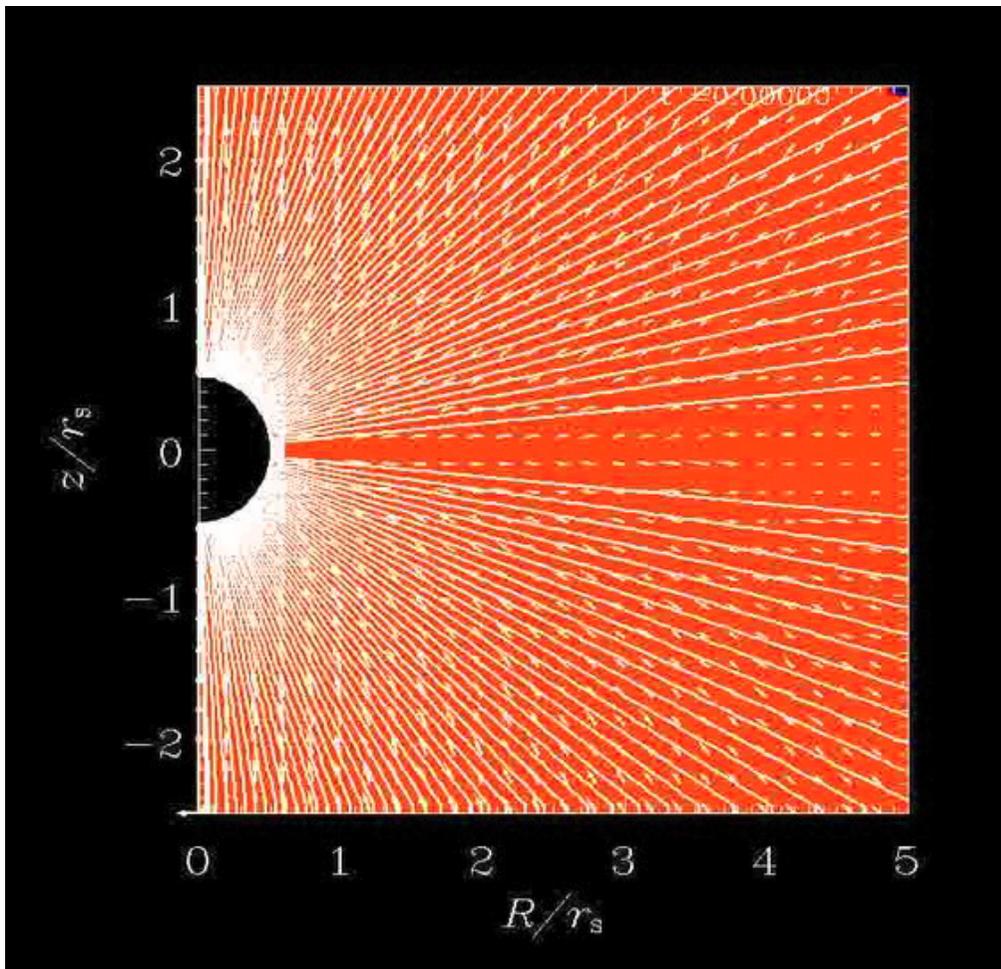
Relativistic magnetic  
reconnection

# Magnetic reconnection in split-monopole magnetic field around black holes

$$\eta = 0.01 r_s$$

- Rapidly rotating (Kerr) black hole case:

$$a = 0.99995$$



Color:  $B_\phi / \rho^{1/2}$

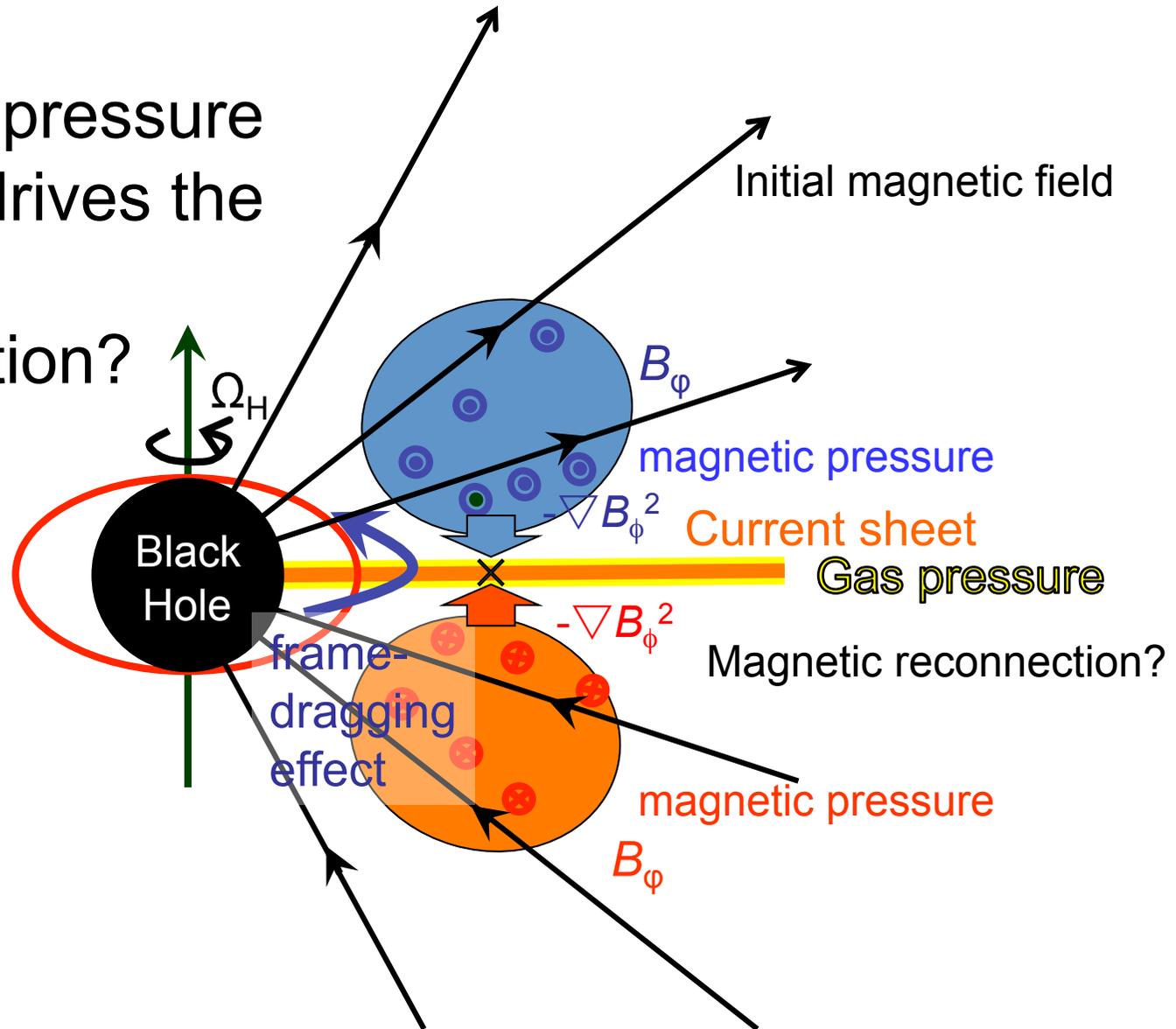
Lines: magnetic field

Arrows: velocity

What mechanism causes the magnetic reconnection in the ergosphere?

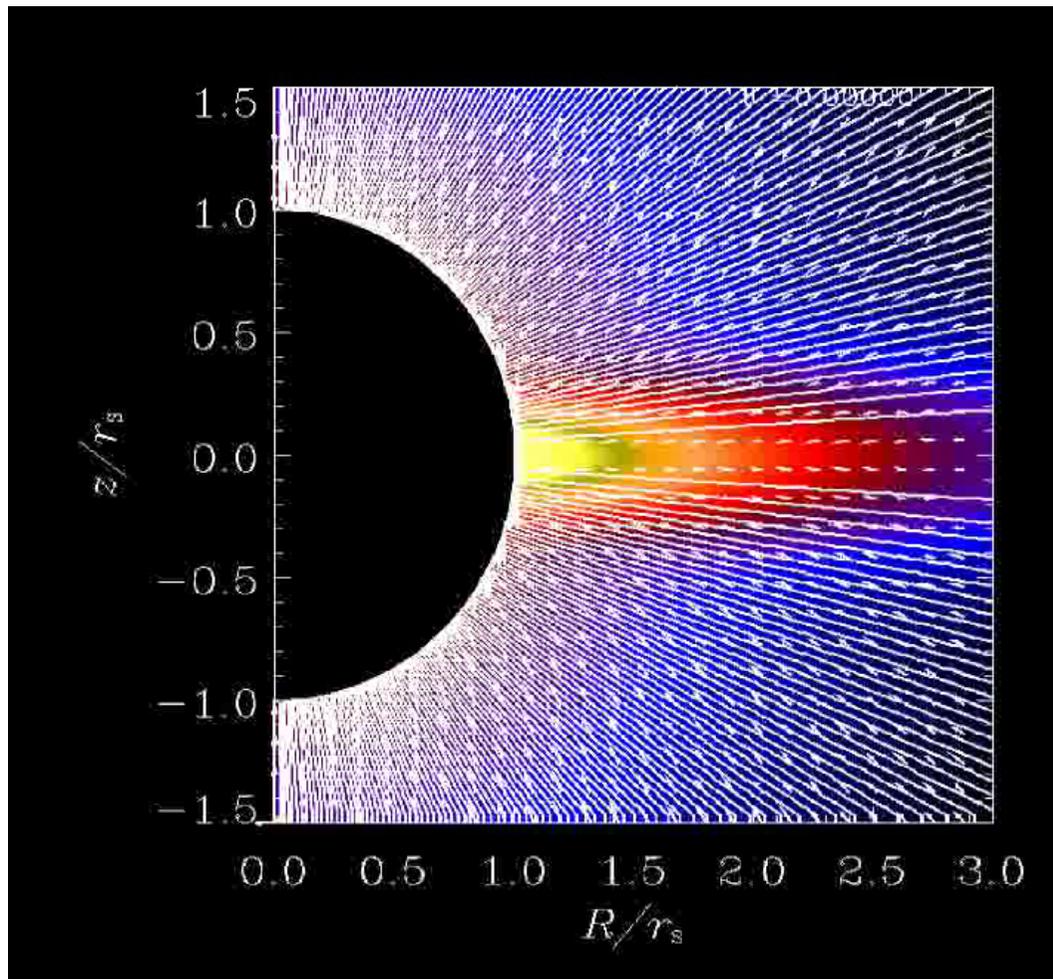
# A expected mechanism of magnetic reconnection around Kerr black hole

Magnetic pressure gradient drives the magnetic reconnection?



# Magnetic reconnection in split-monopole magnetic field around black holes

- Schwarzschild black hole case:



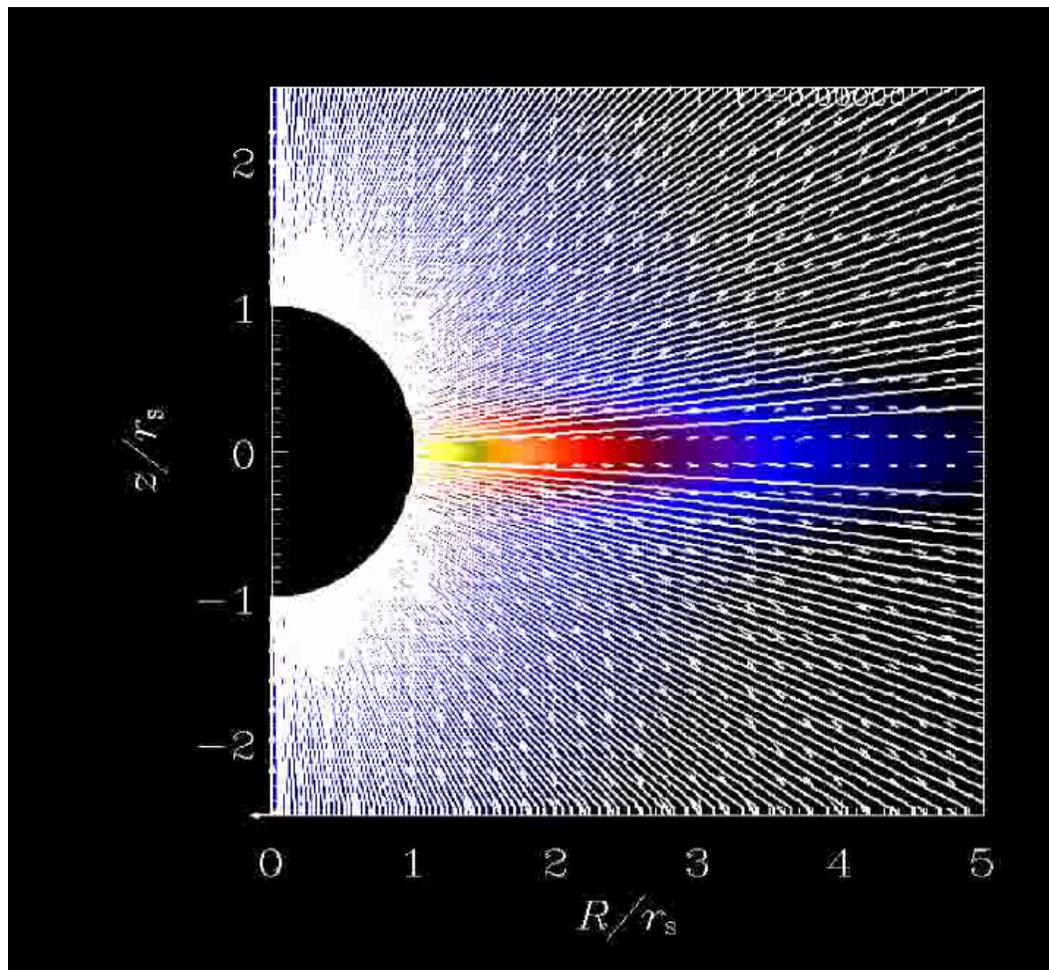
$$a = 0$$

$$\eta = 0.01r_s$$

Color: pressure  
Lines: magnetic field  
Arrows: velocity

# Magnetic reconnection in split-monopole magnetic field around black holes

- Schwarzschild black hole case:



$$a = 0$$

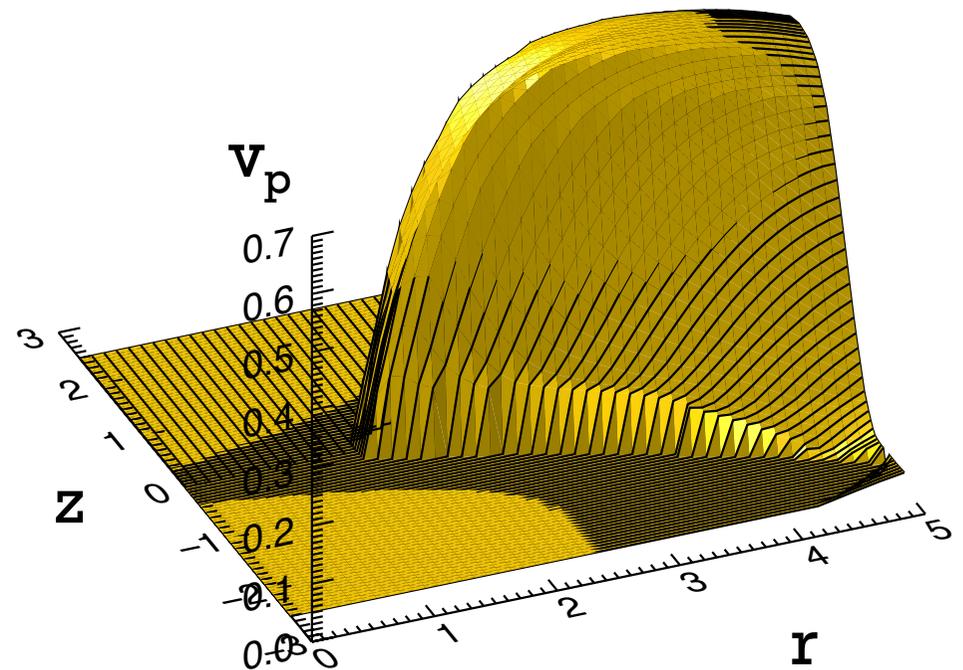
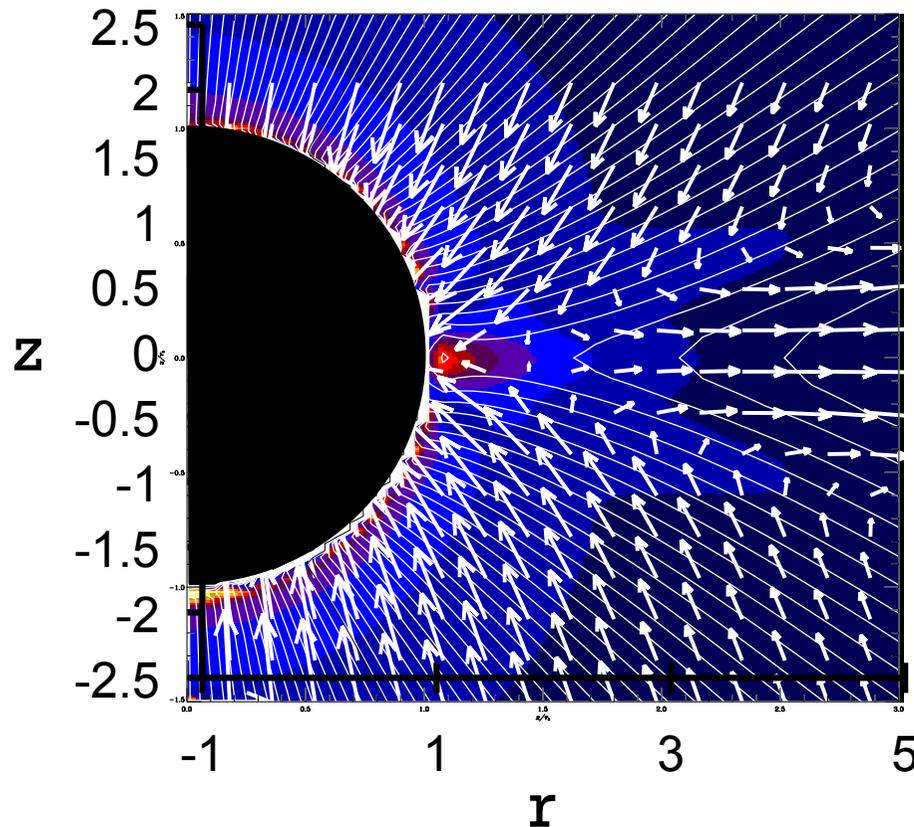
$$\eta = 0.01r_s$$

Color: pressure  
Lines: magnetic field  
Arrows: velocity

# Magnetic reconnection in split-monopole magnetic field around black holes

- Schwarzschild black hole case:

arrow: velocity, color: density



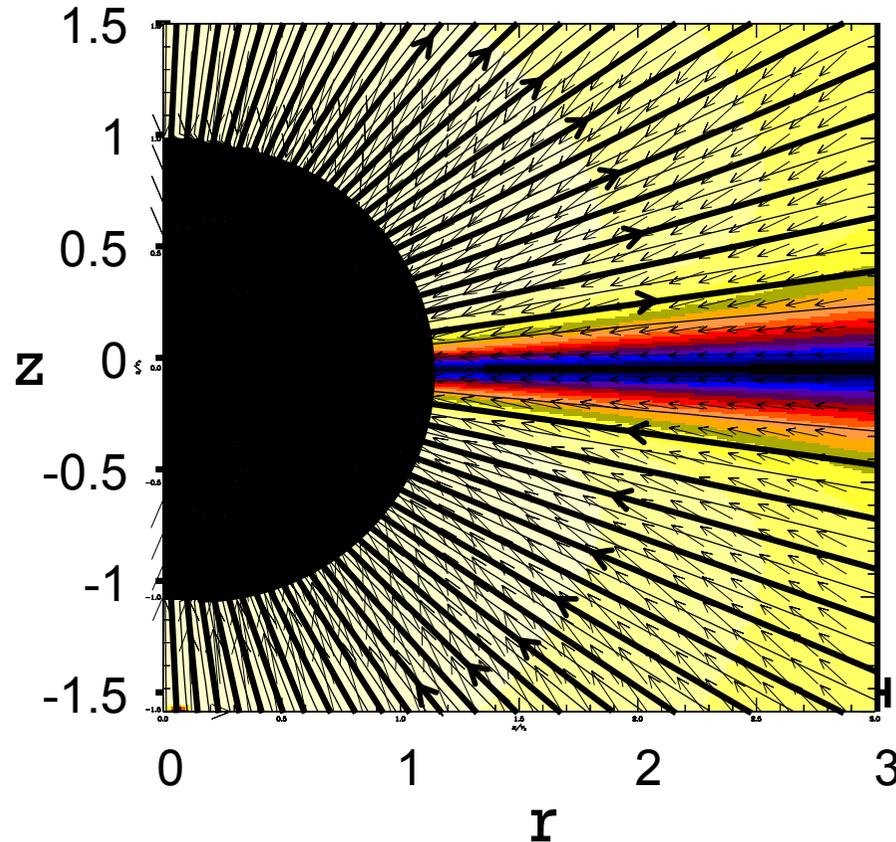
$$v_{p, \max} = 0.62c$$

# Initial condition of magnetic field

$$a = 0$$

Relativistic Alfvén velocity

$t = 0$



$$v_A = \frac{B}{\sqrt{\rho c^2 + \frac{\Gamma}{\Gamma-1} p + B^2}} c$$

Relativistic magnetic reconnection

Split-monopole magnetic field

# Magnetic reconnection rate

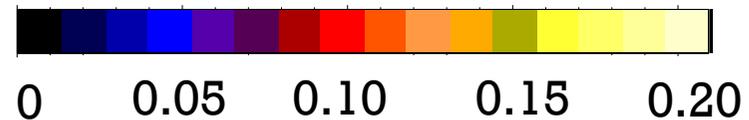
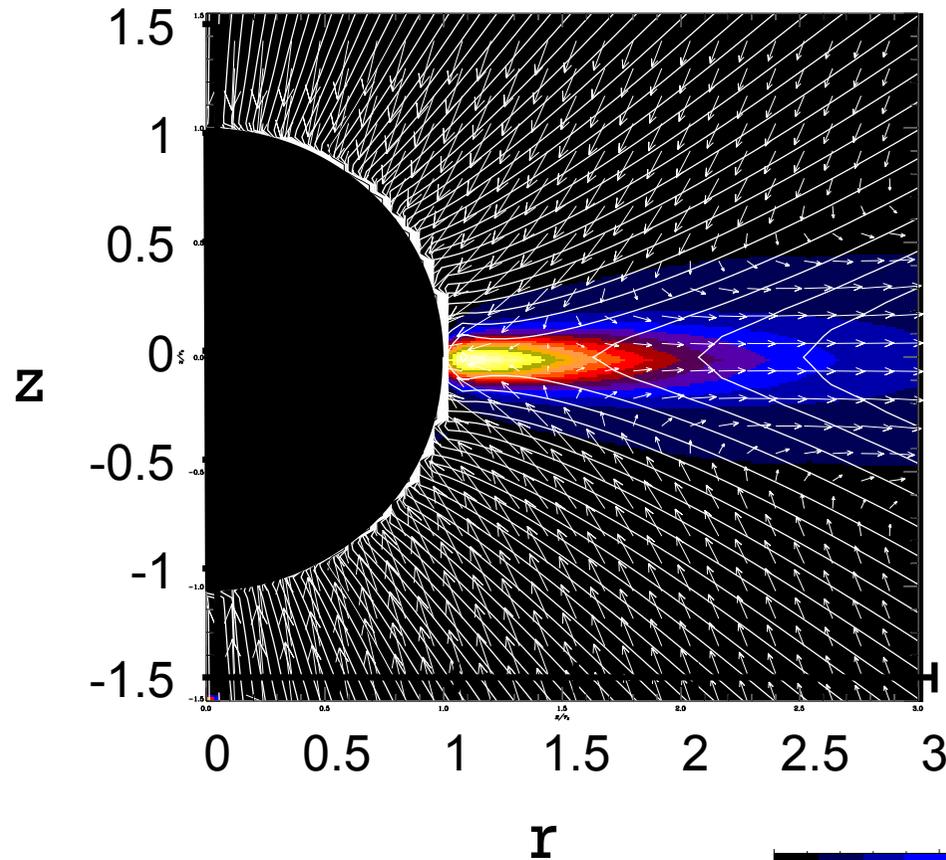
$$a = 0$$

Lapse function

$$\alpha E_\phi = \alpha \eta J_\phi \circ \circ \circ$$

General relativistic  
Faraday's law

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times [\alpha(\mathbf{E} - c\boldsymbol{\beta} \times \mathbf{B})]$$



# Evaluation of tearing instability

Maximum growth rate of tearing instability:

$$\omega_{\max} \cong \tau_A^{-1} R^{-1/2} \quad (\text{Non-relativistic})$$

$$\eta = 0.01 r_S, \quad a_{\text{CS}} = 0.1, \quad v_A = 1$$

$$\tau_A = 0.1 \tau_S, \quad R = 10 \quad \leftarrow \text{Thickness of current sheet}$$

$$\omega_{\max} = 3\tau_S^{-1}$$

$$\lambda_{\max} = \frac{2\pi}{k_{\max}} \approx r_S$$

$$\therefore \omega \cong \tau_A^{-1} (a_{\text{CS}} k)^{-2/5} R^{-3/5}, \quad (ka_{\text{CS}})_{\max} \cong R^{-1/4}$$

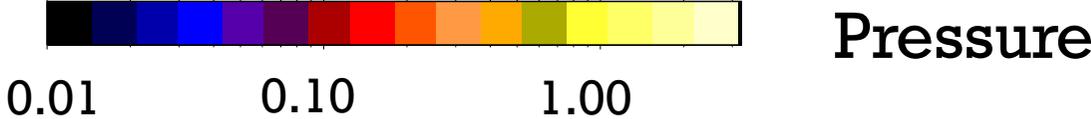
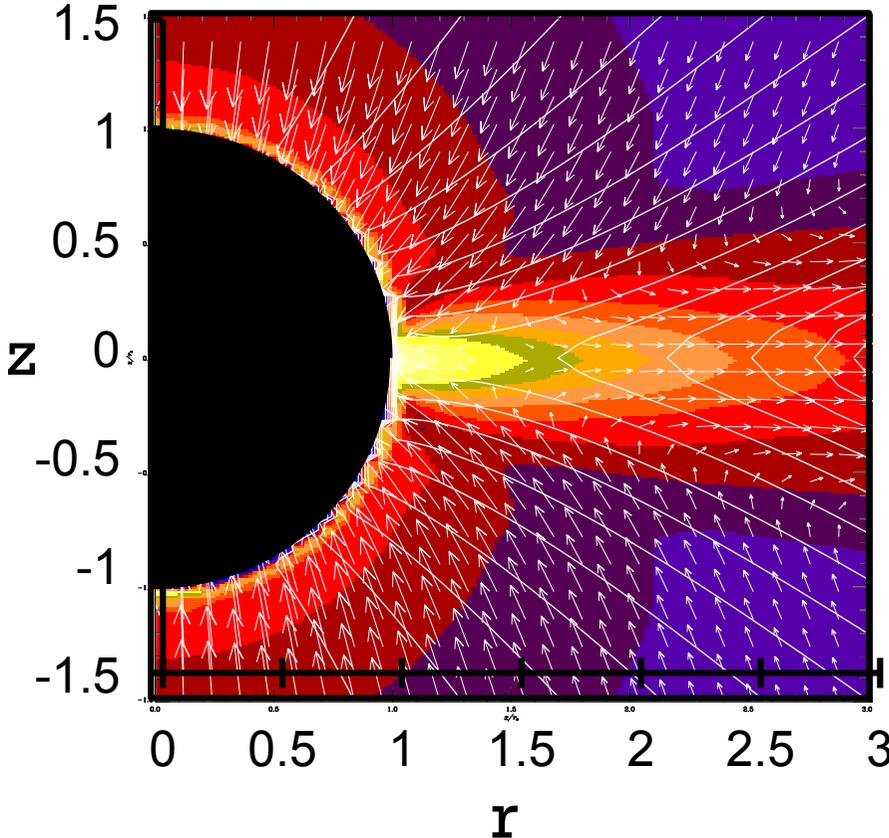
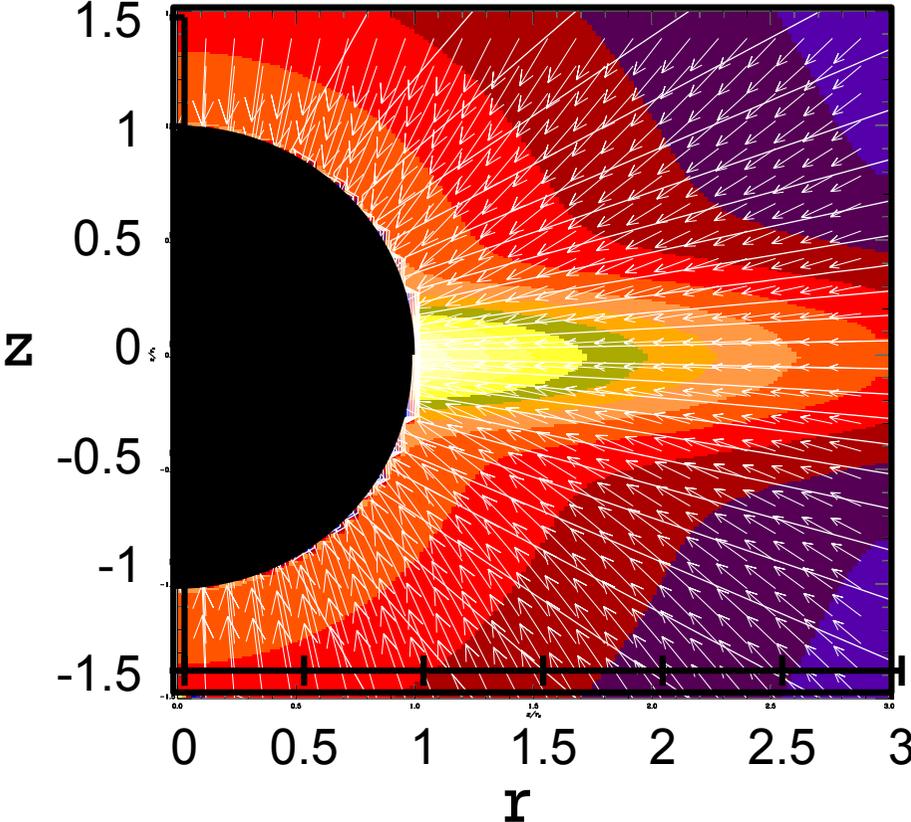
$$\tau_A = \frac{a_{\text{CS}}}{v_A} \quad (\text{Alfven transit time}) \quad R = \frac{a_{\text{CS}} v_A}{\eta} \quad (\text{Magnetic Reynolds number})$$

However, the magnetic reconnection does not depend on the perturbation of the initial velocity at the current sheet. This indicates that the tearing instability has no essential role in the magnetic reconnection.

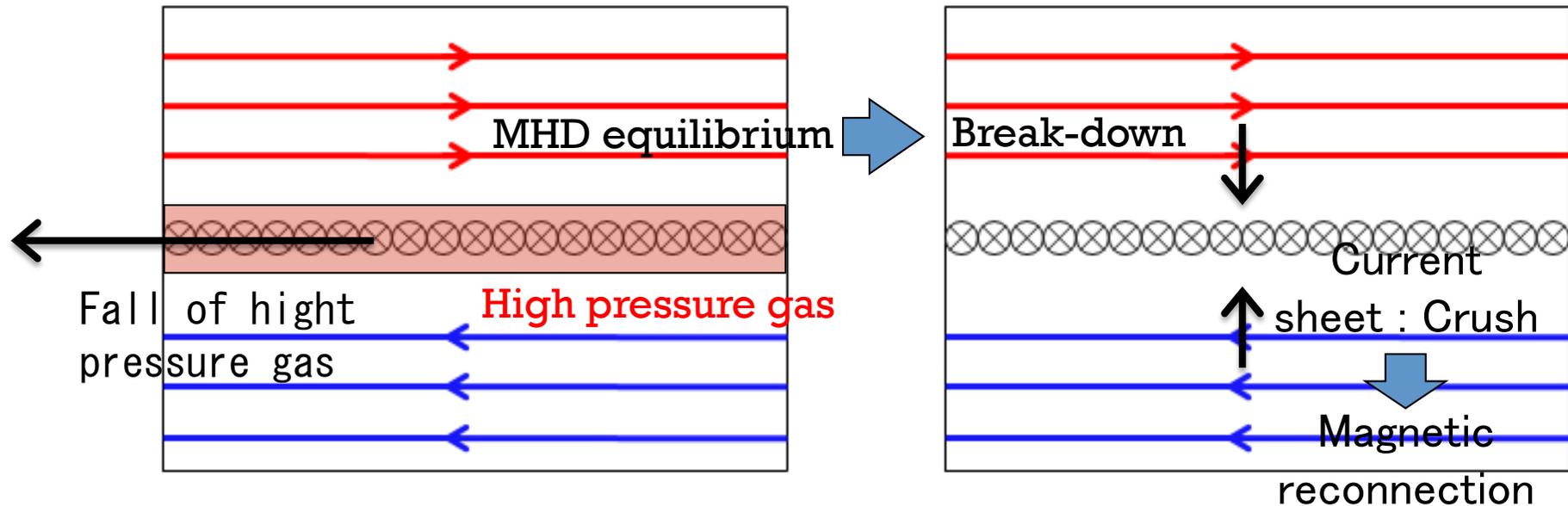
# Analysis of numerical results of magnetic reconnection around Schwarzschild black hole

t=0

t=10

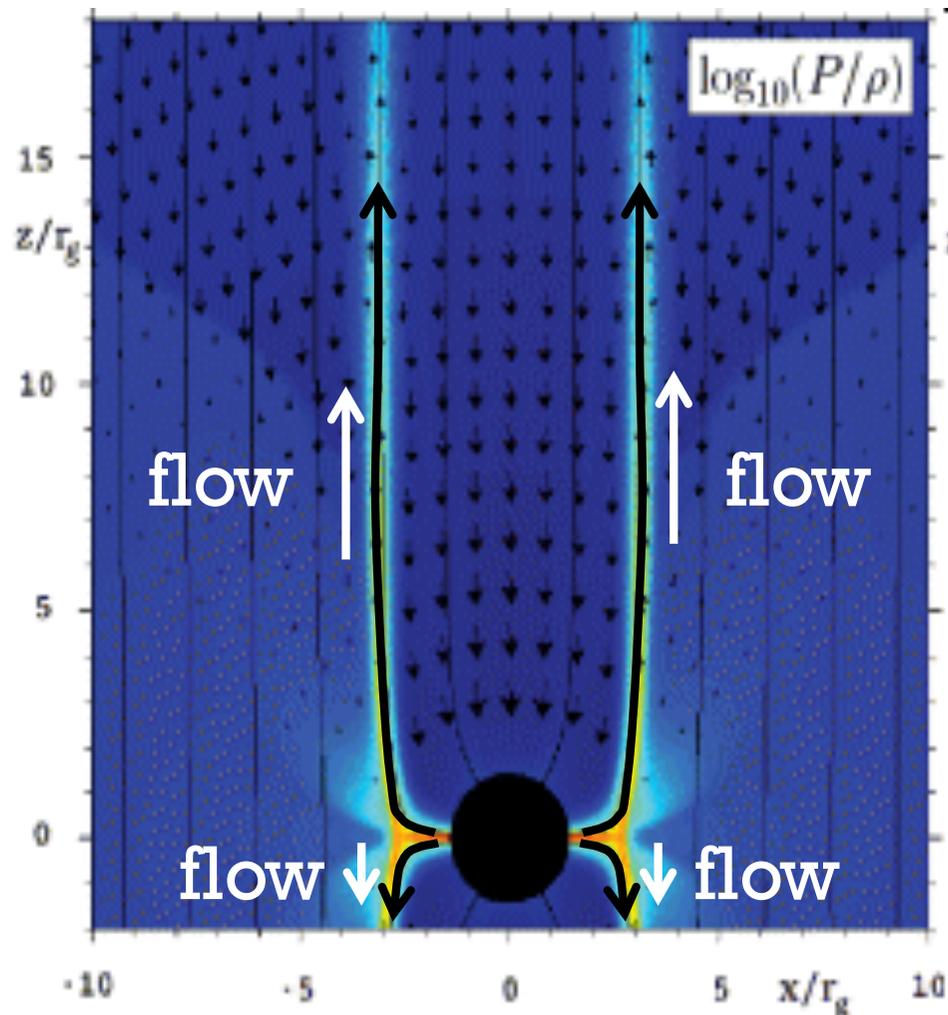


# Preliminary model of the magnetic reconnection in ergosphere: Mechanism of “Daruma-Otoshi” ?



Daruma-otoshi:  
Japanese traditional toy.  
When one piece is hit to  
shut out, the pieces  
upper and lower are  
connected.

# Expected flow from ergosphere driven by magnetic reconnection in initial uniform magnetic field case



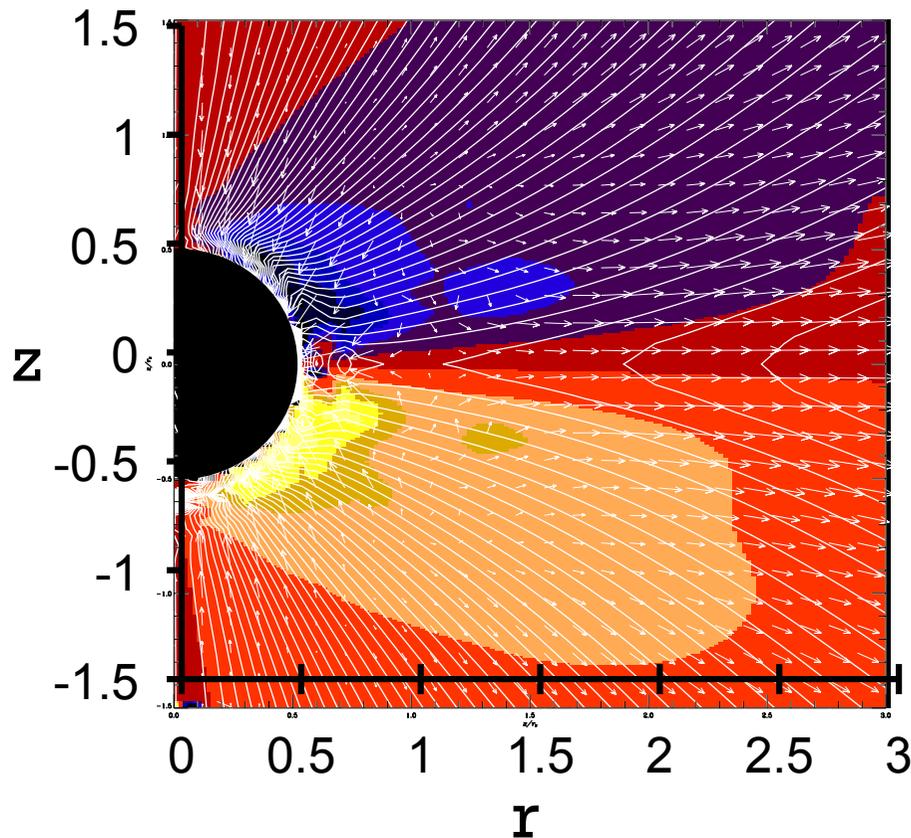
Even in the initial uniform magnetic field case, around the rotating black hole, magnetic reconnection is caused spontaneously and hollow flow from the ergosphere is expected to appear.

Ideal GRMHD simulation of uniformly magnetized plasma around rotating black hole (Komissarov 2005)

# Kerr black hole case

$$a = 0.99995$$

Color: azimuthal component of magnetic field,  $B_\phi$   
Solid lines: magnetic field lines  
Arrows: velocity of plasma



In Kerr black hole case:  
 $v_{p,\max} = 0.71c$  (at  $t = 10 \tau_s$ )

In Schwarzschild black hole case:  
 $v_{p,\max} = 0.63c$  (at  $t = 10 \tau_s$ )

More explosive magnetic reconnection is occurred in Kerr black hole case compared with the case of Schwarzschild black hole case.

# Summary

- I present the basis of resistive GRMHD, with respect to plasma dynamics around black holes, the standard equations, and its causality.
- I performed simulations of resistive GRMHD with split-monopole magnetic field around black holes. The primary results showed spontaneous magnetic reconnection.
- I proposed a model, “daruma-otoshi” mechanism of magnetic reconnection around black holes.
- The numerical results suggest even in initial uniform magnetic field case, the magnetic reconnection is caused spontaneously around the rotating black hole.  
In conclusion, in black hole magnetosphere, it is full of magnetic reconnection!

# Future plans

- Further calculations and analysis of magnetic reconnection around black holes are needed to confirm the mechanism of the magnetic reconnection near the black hole.
- To perform numerical simulations with longer-term, more realistic magnetic field configuration and clarify importance of magnetic reconnection around astrophysical black holes, numerical code with more precise numerical methods, such as HLLD, are required.
- Radiation should be included in GRMHD simulations. I consider “test EM wave” calculation with resistive GRMHD.
- Comparison between magnetic field configurations suggested by our future precise numerical simulations and forthcoming radio observations with high spatial resolution.